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INTERPLAY OF SUPERCONDUCTING AND INSULATING PHASES IN THE METASTABLE  
HIGH-RESISTANCE STATES OF THE  $\text{Ga}_{50}\text{Sb}_{50}$  ALLOY

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Temperature dependences of the resistance  $R(T)$  of the alloy GaSb at different stages of its transformation from the high-pressure metastable metallic (M-) phase into the an insulating (I-) phase are measured. Absolute values of  $R$  at these stages span more than eight orders of magnitude. The superconducting (S-) transition takes place in the low-resistance states at  $T \approx 4.6K$ . It gives way to a steep increase of the resistance, i.e. to an insulating transition at the same temperature in the high-resistance states.

In three dimensions (3D), the quasireentrant superconducting(qrS-) transition usually takes place near the localization threshold [1]. Recently, some experimental indications appeared that qrS-transition can be realized not only in granular material but in the homogeneously disordered material as well [2] or, at least, in materials which cannot be treated as random mixture of conductors and insulators. In this connection, we compared in this paper the  $R(T)$  dependence above the onset of the qrS-transition with that at low temperatures where  $R$  increases with reducing  $T$ .

Our material, the alloy  $Ga_{50}Sb_{50}$ , is one of those alloys which under the high pressure can be obtained in a metastable M-phase and then transformed into I-phase by heating [3]. Dosing the heating, one can get a row of intermediate states and measure the dependence  $R(T)$  in them at low temperatures. The alloy Ga-Sb apparently transforms in accordance with the fractal scheme: the I-phase appears as a fractal structure thinning and entangling the current paths [4]. Indeed, even after the initial resistance of the sample has

been increased by 8 (!) orders of magnitude (estimated starting value of resistivity being  $100 \text{ } \Omega \cdot \text{cm}$ ) we did not reach the state which could be regarded as an I-state judging from formal extrapolation to  $T=0$  of the conductance ( $T$ ).

The first four orders of magnitude in changing  $R$  are not accompanied by any changes of the S-transition. This confirms that only a part of the sample is yet involved into the transformation. The physical properties of those domains where the S-transition preserved remain unchanged but their volume reduces rapidly and their topology becomes more complicate. Further increase of  $R$  leads to appearance of tails in the transition curves first and then to the qrS-transition [1,2]: the resistance fails to become zero and starts to increase with lowering the temperature.

Fig.1 exposes the experimental data. The states are labelled by number  $q$

$$q = \log \left( \frac{R}{R_{in}} \right)_{T=6K}$$

where  $R_{in}$  is the resistance of the sample in the initial state. It can

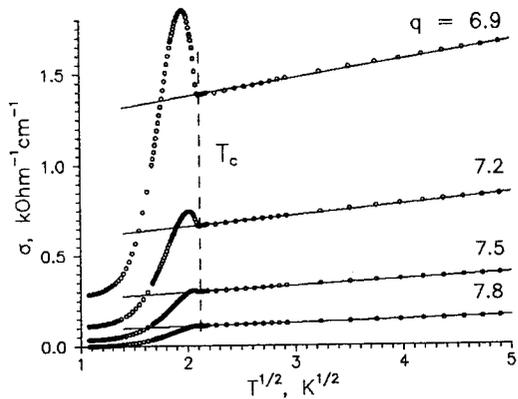


Fig.1

be seen from the curve  $q=7.8$  that in the high-resistance states there is no natural S-response at all: a transition does exist at the temperature  $T_c$  but the resistance at  $T < T_c$  increases instead of decreasing.

It follows from Fig.1 that at  $T > T_c$  and with  $q$  large enough the changes in the conductance are proportional to  $T^{1/2}$ . The straight line

$$\sigma_0(T) = \sigma_0(0) + aT^{1/2}$$

describes the data above  $T_c$ . Below  $T_c$  the difference  $\sigma(T) - \sigma_0(T)$  may change the sign at some temperature  $T^*$ . Let us take  $\sigma_0(0)$  as a parameter of the state and plot  $T_c$  and  $T^*$  vs  $\sigma_0(0)$ . Then we get three domains in the  $(\sigma_0, T)$  plane which can be named M-, S- and I-domains - see Fig.2.

The most natural explanation of the observed phenomenon would suppose that current paths cross a kind of a S-I-S-...-structure with non-Josephson but quasiparticle tunneling currents through I-domains. It is not clear yet whether such approach is valid. Maybe we need quantum

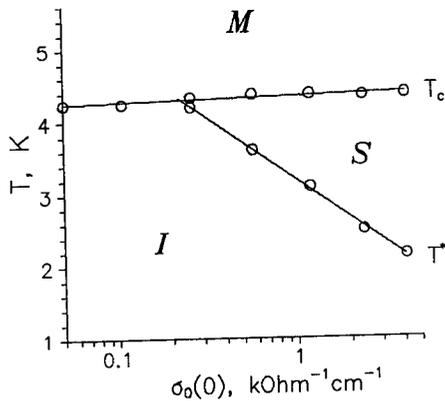


Fig.2

approach instead, treating the system as a whole and involving such ideas as spin density waves, paired electron crystals [5], or supposing finite magnitude of the order parameter and existence of the Cooper pairs at the I-side of the S-I transition [6].

A similar phenomenon - a crossover between S- and I-transitions depending on the conditions of the experiment has been seen in quasi-1D conductor TaSe<sub>3</sub> [7]. The similarity is emphasized by the fact that the dimensions of the conducting domains on a fractal structure may turn to be rather low.

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