Local noise in a diffusive conductor

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The control and measurement of local non-equilibrium configurations is of utmost importance in applications on energy harvesting, thermoelectrics and heat management in nano-electronics. This challenging task can be achieved with the help of various local probes, prominent examples including superconducting or quantum dot based tunnel junctions, classical and quantum resistors, and Raman thermography. Beyond time-averaged properties, valuable information can also be gained from spontaneous fluctuations of current (noise). From these perspective, however, a fundamental constraint is set by current conservation, which makes noise a characteristic of the whole conductor, rather than some part of it. Here we demonstrate how to remove this obstacle and pick up a local noise temperature of a current biased diffusive conductor with the help of a miniature noise probe. This approach is virtually noninvasive for the electronic energy distributions and extends primary local measurements towards strongly non-equilibrium regimes.

Charge transport in electronic conductors1–9 is stochastic by origin, as such is the motion of each individual carrier10. Mean square current fluctuation (current noise) contains valuable information about quasiparticle charge11, transport statistics12–14 and is particularly useful for Johnson-Nyquist15 or non-equilibrium16–20 noise thermometry and even for dynamic measurements of the inelastic relaxation rates21. Microscopically, the noise originates from the random fluctuations in the occupation numbers of the electronic quantum states. Altogether, such fluctuations determine the local noise temperature of the electronic system and give rise to a spontaneous current fluctuation . This fluctuation represents a random current source acting on the conductor at a given location. While in the ballistic regime or in quantum confined conductors is preserved along the device, in general such random sources are mutually uncorrelated at distances exceeding the carrier mean-free path. Current conservation, however, dictates that the actually measurable fluctuation is the same within any cross-section of the conductor and averaged over all sources22, . Here is its length. This fundamental constraint makes conventional noise measurements sample-averaged and hides the information about the local noise temperature.

To overcome this constraint, we propose a concept of a noise sensor depicted in Fig. 1a. The sensor, represented by a thin rod, is brought into contact with a current biased conductor, in which a spatially inhomogeneous non-equilibrium state is indicated by a color gradient. The non-equilibrium electronic energy distribution at the contact position is picked up by the sensor and gives rise to measurable excess noise at a zero net current. A proper non-invasive noise probe has to minimize the disturbance of the electronic energy distributions via the heat leakage. In addition, its generated voltage noise, which contains the desired local information, has to dominate over the spurious voltage noise of the conductor itself. These criteria are naturally fulfilled provided the sensor resistance, , is much higher than that of the conductor, , which can be realized with numerous material combinations and, e.g., in scanning tunneling microscopy23–25. Here, we probe the local noise of a micro-sized constriction in a macroscopic normal metal conductor with the help of a diffusive InAs nanowire (NW), which provides good ohmic contact and is characterized by a convenient resistance ratio of . The spatial resolution in this case is limited by the diameter of the NW. We succeed to observe the shot noise behavior and the lack of electron-phonon relaxation within the constriction, an impossible task for sample-averaging noise approaches. Our approach extends local primary measurements towards a few 10 K temperature range, which is not accessible for spectral sensitive approaches1–3.

The SEM image of the actual device used for local noise measurements is shown in Fig. 1b. In the middle, an n-doped InAs NW (≈70 nm in diameter and 2 μm in length), grown by Au-assisted chemical beam epitaxy26 and drop casted on a SiO2/Si substrate, is visible as a light-gray thin rod. The rest shows a Ti/Au metallic nanostructure around the NW, fabricated with the e-beam lithography. Two ohmic contacts between the NW and metal layers

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are marked as test end (red circle) and cold end (blue circle). The ohmic contacts are realized by means of the two bar-shaped constrictions, each of them further connected to four metallic terminals. These terminals are designed to be broader and twice thicker than the constrictions, to minimize their resistive heating during non-equilibrium driving and improve a thermal coupling to the substrate. In between, the device is additionally equipped with six plunger gates, which are kept grounded throughout the experiment. We control the non-equilibrium noise at the test end position by varying the driving (heating) current $I_H$ which flows through the constriction $H$ on the left hand side. $I_H$ can be switched between two current paths: from the terminals 1 or 2 towards the grounded terminal, referred to as $\perp$ in the following. On the cold end, all the terminals apart from $N$ are kept floating throughout the experiment.

**Results**

**Shot noise calibration.** The performance of the noise probe depends on the way the non-equilibrium fluctuations penetrate into the NW, as determined by the interplay of diffusion and inelastic scattering of the charge carriers. Compared to the elastic transport regime and depending on the nature of the inelasticity, the noise can be either enhanced, as for the electron-electron (e-e) collisions, or suppressed, as for the electron-phonon (e-ph) relaxation. As shown below, none of these effects is relevant to our NWs, which is an ideal case for the noise probe. To understand this better, we perform a primary calibration, biasing the NW with a transport current ($I_{NW}$) via terminal $N$ and measuring the induced shot noise, see Fig. 1b. In this experiment, the constrictions remain in thermal equilibrium at a given bath temperature $T_0$. The result of such calibration is shown in Fig. 2, where we plot the dependence of the measured spectral density ($S_I$) of the current noise on $I_{NW}$ at different $T_0$. At increasing transport current, $S_I$ crosses over from the equilibrium Johnson-Nyquist value of $4k_BT_0/R$ to the closely linear dependence $S_I \approx 2eF|I_{NW}|$, where $F$, $k_B$, and $e$ are, respectively, the Fano factor, the Boltzmann constant, and the elementary charge and $R$ is the linear response resistance of the NW. This observation evidences the shot noise behavior with the Fano factor of $F \approx 0.3$ (dashed line), which is remarkably close to the universal value of $1/3$ in metallic diffusive conductors with negligible inelastic...
scattering\(^{22,28}\). These findings are independent of the device choice, see the inset of Fig. 2, and \( T_0 \) and persist up to \( |I_{\text{NW}}| \approx 1 \, \mu \text{A} \), i.e. the voltage drop of more than 10 mV across the NW. We conclude that the charge transport in our InAs NWs occurs via elastic diffusion, whereas the e-ph relaxation\(^{22}\) and the e-e collisions\(^{27}\) can be safely neglected for carrier excess energies below 10 meV in respect to the Fermi level. Remarkably, this observation is even more favorable, than ab initio estimates of the inelastic scattering rates (Supplemental Material).

The fact that the NW behaves as an elastic diffusive conductor facilitates the calibration of the probe sensitivity to the noise at its test end. The randomness of the carrier diffusion process inside the NW gives rise to a complete loss of correlations between the average occupancy of the quantum state and the direction of the carrier momentum. Still, \( n_k \) retains the dependence on the carrier energy \( \varepsilon \) characterized by the energy distribution function \( f_\varepsilon \equiv f_\varepsilon (x) \). For elastic diffusion, \( f_\varepsilon \) obeys the Laplace’s equation\(^{22}\), that is

\[
\frac{\partial^2 f_\varepsilon}{\partial \varepsilon^2} = 0
\]

with the boundary conditions \( f_\varepsilon (0) \) and \( f_\varepsilon (L) \) set, respectively, by the energy distributions at the test end and the cold end of the NW. Spontaneous current fluctuations in a given cross-section of the NW can be conveniently described in terms of the local noise temperature, which is defined as\(^{22,27}\)

\[
T_N (x) = k_B T_s \int f_\varepsilon (x) (1 - f_\varepsilon (x)) \, d\varepsilon.
\]

These fluctuations sum up\(^{22}\) such that the \( T_N (x) \) averages along the device to give response of the noise sensor:

\[
T_N = \int T_N (x) \frac{dx}{L}, \quad S_I = \frac{4k_B T_s}{R}
\]

where \( T_s \) and \( R \) are, respectively, the measured noise temperature of the NW-sensor and its resistance.

In the regime of shot noise calibration, \( f_\varepsilon (0) \) and \( f_\varepsilon (L) \) represent the equilibrium Fermi distributions with the chemical potential difference determined by the voltage drop across the NW. Here, the eqs (1) and (2) predict the well-known double-step distribution in the middle of the NW, see Fig. 1c, and the universal \( F = 1/3 \) for the shot noise Fano factor\(^{22}\). The regime of local noise sensing, addressed below, is different in two respects: (i) the unknown energy distribution at the test end is non-equilibrium and (ii) no net current flows through the NW, \( I_{\text{NW}} \approx 0 \). Here, the eqs (1) and (2) relate the measured noise of the sensor and the local noise temperature of the conductor under test. In the limit of strong non-equilibrium, \( T_N (0) \gg T_0 \), we arrive simply at \( T_N \approx \alpha T_N \) (0), that is the sensor directly measures the noise temperature of the conductor at the position of their mutual contact. As discussed above, elastic diffusive transport in our NWs persists up to at least 10 meV excitation energies, making our local noise thermometry well applicable within a few 10 K temperature range. A pre-factor \( \alpha \) is slightly sensitive\(^{28}\) to the shape of \( f_\varepsilon (0) \), e.g., \( \alpha = \frac{2}{3} \) for a double-step shaped distribution, similar to that sketched in Fig. 1c, and \( \alpha = 1 + \ln 2 \approx 0.56 \) for an equilibrium Fermi-Dirac distribution. For definiteness, the a-priori unknown energy distribution at the test end is assumed below to be of equilibrium type, which is not important for any of our conclusions.
Sensing the local noise. Next we describe the main result of our paper—the proof of concept of the local noise sensor. In this experiment, the noise of the unbiased NW-sensor is measured in response to the driving current $I_D$ sent via the constriction or bypassing it, respectively, in biasing configurations $1 \rightarrow \perp$ and $2 \rightarrow \perp$, see the micrograph in Fig. 1b.

Figure 3 shows the measured noise temperature of the sensor $T_S$ in dependence on $I_D$ at $4.2$ K for the device NW1 (solid symbols) in the biasing configuration $1 \rightarrow \perp$. By definition, the absolute values of $T_S$ translate into a measured voltage noise as $\delta V^2 \approx 4k_B T_S R \Delta f$, where $\Delta f$ is the measurement bandwidth and $R \sim 10$ k$\Omega$ is the resistance of the NW, which is much higher than that of the metallic terminals connecting it to the ground, $r_M \sim 20$ $\Omega$.

On the other hand, $\delta V^2$ consists of two contributions, the NW noise $\delta V^2_{NW}$ and the terminals noise $\delta V^2_M$. Obviously, $\delta V^2_M \ll \delta V^2_{NW}$, since the noise temperature of the metallic terminals, $T_M$, is the same order of magnitude as the temperature at the test end, $T_n(0)$, and $T_n(0) \sim T_S$ as we derived above. Note also, that assuming $\delta V^2_M \sim \delta V^2$ one would arrive at an unreasonable estimate of $T_M \sim (R/r_M) T_S \sim 1000$ K. This remarkable feature of Fig. 3 demonstrates that the voltage noise of the metallic terminals can be safely neglected and $T_S \approx T_{NW}$. Therefore, in agreement with the eqs (1) and (2), the measured noise is generated in the NW in response to a local non-equilibrium carrier distribution penetrating via its test end. In comparison, the noise response is considerably weaker in a biasing configuration $2 \rightarrow \perp$, for which $I_D$ bypasses the constriction, shown by open circles in Fig. 3. This additionally verifies that the measured noise is a local characteristic.

$T_S$ demonstrates a parabolic increase near the origin and an almost linear dependence at large $I_D$. The tendency towards a linear dependence of $T_S$ on the drive current for the upper trace in Fig. 3 is consistent with the shot noise dominated regime of the current biased constriction, which takes place at inefficient electron-phonon (e-ph) energy relaxation. In this regime, under the assumption of local equilibrium, one expects a standard spatial temperature distribution along the constriction, determined by a balance of the Joule heating and the Wiedemann-Franz thermal conduction\(^2\). The maximum local temperature is attained at the test end and equals

$T_N(0) = \sqrt{3e^2/(4\pi k_B^2)} V^2 + T_c^2$, where $V = R I_D$ is the voltage drop on the constriction and $T_c$ is the temperature in the metallic leads. The solid line in Fig. 3 corresponds to a fit parameter $r = 3.7$ $\Omega$ and $T_c = T_M$, i.e., negligible Joule self-heating of the leads. Still, in this experiment the bath temperature is too high to rule out the impact of e-ph energy relaxation in the constriction. For comparison, we performed model calculations of spatial temperature distribution in the constriction and the leads for the two biasing configurations and with the Wiedemann-Franz heat conduction and the e-ph relaxation taken into account, see the insets of Fig. 3. Reasonable fits of the sensor response could be obtained for the e-ph cooling power density of $\Sigma (T^3 - T_0^3)$, with $\Sigma = 0.1$ nW $\mu$m$^{-3}$ K$^{-2}$, see the dashed lines in Fig. 3.

The inefficiency of the e-ph relaxation becomes evident at a lower bath temperature. In Fig. 4 we plot the results of the local noise sensing at $T_B = 0.5$ K in three devices, all in the biasing configuration $1 \rightarrow \perp$. We observe a closely linear dependence on $I_D$ over an order of magnitude increase of $T_S$ (symbols). This behavior is qualitatively different from a sub-linear dependence of $T_S$ on $I_D$, which is expected in presence of e-ph relaxation, see the dashed line fit, obtained with the same parameter $\Sigma$ as the fits of Fig. 3. By contrast, the data is perfectly consistent with the linear shot noise behavior, see the solid line fit, which corresponds to the shot noise of the constriction.
with $r = 3.6 \Omega$. Similar to the case of $T_0 = 4.2 \text{ K}$, this $r$ value is somewhat higher than typically measured four-terminal resistances of $2.8 \div 2.9 \Omega$ in our devices, which is a reasonable discrepancy in view of the constriction geometry, see Fig. 1b. The shot noise fit of Fig. 4 fails to describe a little downward kink at $|I_H| \approx 0.4 \text{ mA}$, more or less pronounced in all our devices. The origin of this kink is not understood and might be related to peculiarities in the convection$^{30}$ or even boiling of the surrounding liquid $^3$He. The observation of shot noise behavior in Fig. 4 demonstrates a success of our local noise measurement approach. Our short constrictions are connected via macroscopic leads, each with a typical resistance of $\sim 20 \Omega$. Thus, in conventional two-terminal noise measurements, even a minor increase of the Johnson-Nyquist noise of the leads, owing to their Joule self-heating, would completely mask the constriction shot noise$^{31}$.

The shot noise behavior in short constrictions is a consequence of hot electron out-diffusion in the leads, which occurs faster than the time-scale of e-ph relaxation, $\tau_{\text{diff}} \ll \tau_{\text{e-ph}} \sim T^{-3}$. An estimate based on a free-electron model$^{32}$ is roughly consistent with the experiment, giving $\tau_{\text{diff}} \sim \tau_{\text{e-ph}} \sim 150 \text{ ps at } T = 10 \text{ K}$. In longer constrictions the diffusion time increases as $\tau_{\text{diff}} \sim L^2$, and we expect the e-ph relaxation to take over. This is demonstrated in a second generation devices, in which a NW is divided in two sections, each serving as a local noise probe, see the inset of Fig. 5. The local noise signals, which arise from biasing either a short constriction H1 or a long meander-shaped heater H2, are picked up at the contact N. These signals are plotted in the body of Fig. 5, symbols. The noise of the short constriction H1, circles has a linear dependence on current, which is perfectly consistent with the shot noise fit for $r = 3.1 \Omega$, see the lower solid line. Remarkably, the shot noise behavior is shown to persist here up to $I_H = 3 \text{ mA}$, corresponding to local temperatures of $T_S(0) \approx 30 \text{ K}$ in the center of the constriction. By contrast, the noise of the long heater H2, squares, is much smaller than its expected shot noise value, upper solid line, and exhibits a sub-linear current dependence, evidencing a dominant role of the e-ph relaxation. Yet, in this case, the measured $T_S$ is considerably higher than the estimate obtained under assumptions of the dominant e-ph relaxation, $P_J = \Sigma(T^5 - T_0^5)$, where $P_J$ is the Joule power dissipated per unit volume of the heater H2, see the dashed line in Fig. 5. This indicates that e-ph cooling is bottlenecked by Kapitza resistance and/or poor substrate heat conduction. From these data, we estimate that in our short constrictions the e-ph relaxation dissipates just a few percent of the total Joule heat, explaining the observed shot noise behavior.

**Discussion**

In summary, we realized a local noise measurement in a current biased metallic conductor using InAs nanowire as a miniature noise probe. This approach enabled us to measure a shot noise of a micro-sized constriction connected to macroscopic resistive leads, an impossible task for conventional noise approaches. Our experiment represents a conceptually simple primary local measurement, not limited to very low temperatures and small lead resistances and compatible with various material combinations. As such, we envision, that local noise probing will turn useful for non-invasive spatially resolved studies in non-equilibrium nano-circuits.
Methods

Au-assisted Se doped InAs NWs are grown by chemical beam epitaxy on an InAs(111) B substrate. The carrier density of the InAs NWs derived by field effect measurements is about $1 \times 10^{18} \text{cm}^{-3}$. Typical ohmic contact resistance in our devices is below $100 \Omega$, whereas the NW resistance is about $10 \text{k} \Omega$ per micrometer. We performed the measurements in two $^3$He inserts, with the samples immersed in liquid (at $T = 0.5 \text{ K}$) or in gas (at $T = 4.2 \text{ K}$) and placed vertically face down. The shot noise spectral density was measured using home-made low-temperature amplifiers (LTamp) with a voltage gain of about $10 \text{ dB}$, input current noise of $\sim 10^{-27} \text{A}^2/\text{Hz}$ and dissipated power of $\sim 200 \mu \text{W}$. We used a resonant tank circuit at the input of the LTamp, see the sketch in Fig. 1b, with a ground bypass capacitance of a coaxial cable and contact pads $\sim 40 \text{ pF}$, a hand-wound inductance of $\sim 6 \mu \text{H}$ and a load resistance of $10 \text{k} \Omega$. The output of the LTamp was fed into the low noise 75 dB total voltage gain room temperature amplification stage followed by a hand-made analogue filter and a power detector. The setup has a bandwidth of $\Delta f \sim 0.5 \text{ MHz}$ around a center frequency of $\approx 10 \text{ MHz}$. A calibration was achieved by means of equilibrium Johnson-Nyquist noise thermometry. For this purpose we used a commercial pHEMT transistor connected in parallel with the device, that was depleted otherwise. All transport measurements were performed with the help of a two-terminal or four-terminal lock-in resistance measurement.

References


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Author Contributions

D.E., F.R. and M.R. fabricated the devices, E.S.T. performed experiments and prepared the figures. E.S.T., D.V.S. and V.S.K. analyzed the data. L.S., S.R. and V.S.K. have jointly supervised the work and written the manuscript. All authors reviewed the manuscript.

Additional Information

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