Competition and coexistence of antiferromagnetism and superconductivity in $\text{R}$$\text{Ba}_2$$\text{Cu}_3$$\text{O}_{6+\delta}$ ($\text{R}=$ Lu, Y) single crystals

A. N. Lavrov,$^1$ L. P. Kozeeva,$^1$ M. R. Trunin,$^2$ and V. N. Zverev$^2$

$^1$Institute of Inorganic Chemistry, Lavrentyeva-3, Novosibirsk 630090, Russia
$^2$Institute of Solid State Physics, Institutskaya-2, Chernogolovka 142432, Moscow Region, Russia

Received 29 April 2009; published 22 June 2009

We use $c$-axis resistivity and magnetoresistance measurements to study the interplay between antiferromagnetic (AF) and superconducting (SC) ordering in underdoped $\text{R}$$\text{Ba}_2$$\text{Cu}_3$$\text{O}_{6+\delta}$ ($\text{R}=$ Lu, Y) single crystals. Both orders are found to emerge from an anisotropic three-dimensional metallic state upon which antiferromagnetism opposes superconductivity by driving the doped holes toward localization. Despite the competition, the superconductivity sets in before the AF order is completely destroyed and coexists with latter in a certain range of hole doping. We find also that strong magnetic fields affect the AF-SC interplay by both suppressing the superconductivity and stabilizing the Néel order.

DOI: 10.1103/PhysRevB.79.214523  PACS number(s): 74.25.Fy, 74.25.Dw, 75.30.Kz, 75.47.--m

I. INTRODUCTION

In high-$T_c$ cuprates, the conducting state develops gradually as the parent antiferromagnetic (AF) insulator is doped with charge carriers. At a certain point of this transformation the superconductivity (SC) sets in, which happens long before such remnants of the AF insulating state as AF correlations and a pseudogap in the electron density of states are completely wiped out. Despite the decades of intensive research, however, detail mechanisms of the insulator-metal transformation and, in particular, the role of AF correlations therein remain unclear. From the experimental point of view, the difficulty stems from the very strong exchange interaction in CuO$_2$ planes, $J\sim 0.1$ eV, which makes the AF correlations robust to high temperatures and magnetic fields and leaves no means of "switching" them off. For compositions on either side of the insulator-superconductor transformation, two-dimensional (2D) AF correlations nucleate in CuO$_2$ planes already at high-temperatures $T\sim J/k_B$. As the temperature decreases, the spin correlations smoothly develop until—depending on the doping level—establishing a long-range AF order at the Néel temperature $T_N$, forming a static or fluctuating "stripe" structure, or freezing into a "cluster spin glass." This evolution being spread over a wide temperature range makes it hard to clarify whether and how the spin textures emerging in CuO$_2$ planes affect the charge motion.

In contrast to the gradual development of 2D AF correlations within the CuO$_2$ planes, the spin ordering along the $c$ axis—if it takes place at all—occurs in a dramatically compressed temperature interval. Owing to the strong anisotropy of exchange interactions, the $c$-axis spin correlation length $\xi_c^{AF}$ stays shorter than the interplane distance until the three-dimensional (3D) ordering temperature $T_N$ is closely approached. Only in the vicinity of $T_N$, $\xi_c^{AF}$ reaches the unit-cell size, causing a 2D–3D crossover in AF correlations and then diverges sharply. In this paper, we report a $c$-axis magnetoresistance (MR) study of $\text{Lu}$$\text{Ba}_2$$\text{Cu}_3$$\text{O}_{6+\delta}$ (Lu-123) and $\text{Y}$$\text{Ba}_2$$\text{Cu}_3$$\text{O}_{6+\delta}$ (Y-123) single crystals with doping levels spanning the range of the AF-SC transformation. Sharp AF transitions observed in crystals even with low $T_N=20$ K have allowed us to map precisely the phase diagram and to single out the impact of spin ordering on the $c$-axis charge transport. We find that in the doping region of the AF-SC transformation the (Lu,Y)-123 crystals exhibit a 3D-metallic conduction until the SC or AF ordering intervenes at low temperatures. The Néel ordering appears to compete with SC by strongly inhibiting the charge motion, yet nevertheless the two orders coexist in a certain range of hole doping. Interestingly, the magnetic field is found to interfere with both orders suppressing the superconductivity and stabilizing the static AF order.

II. EXPERIMENTAL DETAILS

$\text{R}$$\text{Ba}_2$$\text{Cu}_3$$\text{O}_{6+\delta}$ single crystals with nonmagnetic rare-earth elements $\text{R}=$ Lu and Y were grown by the flux method and their oxygen stoichiometry was tuned to the required level by subsequent high-temperature annealing and quenching. Given the complex distribution of structural defects along the $c$ axis in flux-grown crystals, part of the crystals were cut and polished to select better pieces for measurements. We confirmed also that no oxygen-enriched layer was formed at the crystal surface upon quenching; namely, one of the crystals was dissolved in acid in several steps, each step being preceded with resistivity measurements, yet no considerable change in the SC transition was observed.

After the doping level of $\text{R}$$\text{Ba}_2$$\text{Cu}_3$$\text{O}_{6+\delta}$ crystals was roughly set by changing the oxygen content $x$, the hole density in CuO$_2$ planes $n_h$ was further reversibly tuned using the...
well-studied phenomenon of chain-layer oxygen ordering: heating the crystal above 100–120 °C with subsequent quenching reduced the hole density in CuO$_2$ planes by up to 20%, while room-temperature aging gradually restored it back without any change in stoichiometry.$^{8,9}$ For characterizing the doping state of crystals we have selected the in-plane conductivity $\sigma_{ab}(280$ K). The latter is a good measure of the hole density given that the hole mobility stays almost constant in the doping region under investigation.$^2$ The particular temperature $T=280$ K is taken to be high enough to reduce the impact of impurity scattering, yet sufficiently low to avoid oxygen ordering and thus change in doping on the time scale of measurements. It is worth noting that although the in-plane and the $c$-axis transport measurements were carried out on separate samples, the samples were cut from the same crystal and then annealed and handled together to provide an exact match in their states.

The $c$-axis resistivity $\rho_c(T)$ was measured using the ac four-probe technique with a special care taken to avoid mixing of $\rho_c$ with the in-plane resistivity $\rho_{ab}$. To provide a homogeneous current flow along the $c$ axis, two current contacts were painted to almost completely cover the opposing $ab$ faces of the crystal, while two voltage contacts were placed in small windows reserved in the current ones.$^8$\textsuperscript{10}$ Magnetoresistance measurements were carried out either by sweeping temperature at fixed magnetic fields $H//c$ up to 16.5 T or by sweeping the field at fixed temperature.

III. RESULTS

A. $n_x$-$T$-$H$ phase diagram

In RBA$_2$Cu$_3$O$_{6+x}$, the hole density in CuO$_2$ planes, $n_x$, can be tuned continuously by changing the degree of oxygen ordering (see Sec. II). This provides a number of states with different $n_x$ for a given oxygen content $x$, thereby allowing critical regions of the phase diagram to be studied in detail. We have used this approach to trace how the $c$-axis resistivity of LuBa$_2$Cu$_3$O$_{6+x}$ single crystals evolves with the hole density in the narrow doping region where superconductivity sets in (Fig. 1). As can be seen, the $\rho_c(T)$ curves in this heavily underdoped region demonstrate a complicated behavior with a crossover$^{12}$ at 100–150 K. At temperatures above the crossover, the $c$-axis resistivity smoothly decreases with increasing hole density and no peculiarity shows up in this smooth evolution at the doping level of the SC onset. The low-$T$ resistivity behavior, however, changes remarkably. From the set of zero-field $\rho_c(T)$ curves in Fig. 1 one can immediately infer that, upon cooling, the crystal either turns into the SC state or instead exhibits a steep increase in the resistivity. This steep growth in $\rho_c(T)$ with a characteristic asymmetric negative peak in the derivative (inset in Fig. 1) is a hallmark of the Néel transition in RBA$_2$Cu$_3$O$_{6+x}$.$^8$\textsuperscript{10}$ The evolution of $\rho_c(T)$ curves looks therefore as if AF and SC are competing for carriers and the system falls into one of two possible states, either the insulating AF state or the SC one. It is worth emphasizing that sharp AF transitions are observed down to remarkably low $T_N=20$ K (inset in Fig. 1) instead of the freezing into a spin glass usually expected at such temperatures,$^5$ the superconductivity in Lu-123 thus develops directly from the AF state without any intervening spin-glass region.

One may wonder whether a strong magnetic field will suppress both the AF and SC orders and create some intermediate metallic state. We find, surprisingly, that upon suppressing the superconductivity the field $H//c$ also recovers the steep increase in $\rho_c$ associated with the AF ordering (Fig.

FIG. 1. (Color online) Out-of-plane resistivity, $\rho_c(T)$ of a LuBa$_2$Cu$_3$O$_{6+x}$ single crystal at several hole-doping levels near the AF-SC transformation; the hole doping was tuned by changing the oxygen content $x$ and the degree of oxygen ordering (several states for each $x$). The data were taken at zero and 16.5 T magnetic field applied along the $c$ axis. Inset: An anomaly in the normalized derivative $(d\rho_c/dT)/\rho_c$ associated with the Néel transition; measurements were done after quenching the crystal (●) as well as after room-temperature aging for several hours (▲) and days (○).

FIG. 2. (Color online) Suppression of the superconductivity with the magnetic field in Lu-123 and Y-123 crystals. The unveiled normal-state resistivity exhibits an anomaly associated with the Néel transition (compare with the inset in Fig. 1).
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The competition and coexistence of sharp for transition temperatures reduced by 20–25 times mechanisms underlying the observed conductivity density. A study of other crystals has confirmed that as long as the SC is weak enough to be wiped out with the 16.5 T field, be that in Lu-123 or Y-123 crystals, the obtained normal-state SC is weak enough to be wiped out with the 16.5 T field. Moreover, by aging the crystals at room temperature and measuring the Tc of several crystals systematically driven through the entire doping range, the Tc line at ~15 K, contrasting to another prototype high-Tc cuprate La2−xSrxCuO4 where the long-range AF order is restored long before the SC sets in.1,5

An important issue is the reliability of the presented phase diagram, especially given the controversy remaining in publications.4,5,13,14 Indeed, the observed SC and AF transitions are not very sharp, leaving some room for doubts. Besides, however sharp the transitions might be, a possibility of a peculiar macroscopic inhomogeneity imitating the phase coexistence will still remain. For the oxygen compositions under discussion, we have ruled out one possible source of inhomogeneity associated with a phase separation into oxygen-rich and oxygen-poor domains: the behavior of Lu-123 crystals quenched from high temperatures and thus possessing homogeneous oxygen distribution, showed no qualitative difference from that of crystals aged at room temperature for an ultimate 10-years period.

A static cation inhomogeneity over the crystal as a source of the AF and SC phase coexistence is, however, hard to rule out by conventional means. Fortunately, in R-123 crystals one has an additional tool to check the phase diagram by employing the hole doping though the oxygen ordering.8,9 In fact, by aging the crystals at room temperature and measuring their conductivity σab(280 K), Tc, and TN as a function of time, one can evaluate the slopes of phase boundaries in the diagram shown in Fig. 3. An important point here is that the relative changes in Tc and TN observed upon aging are insensitive to static macroscopic inhomogeneities. The phase-
boundary slopes measured in the described way are presented in Fig. 4 as a function of $T_c$ and $T_N$; this way of the data presentation is selected to avoid any uncertainty related to ill-defined parameters such as the hole density. As expected, the slope of the SC phase boundary, $\Delta T_c/\Delta \rho_{\text{ao}}$ (280 K), gradually increases with decreasing $T_c$ and the superconductivity disappears abruptly at a critical doping. The behavior of the AF boundary is much different; initially it indeed becomes steeper with increasing doping and decreasing $T_N$, but then the slope of the $T_N$ line decreases dramatically tending to zero at zero $T_N$. This flattening unambiguously confirms the intrinsic origin of the AF-order tail entering the SC region in the phase diagram in Fig. 3.

What still remains to be clarified is whether the AF and SC orders actually coexist in their overlapping region on the phase diagram or the AF order is restored by the magnetic field upon destroying the superconductivity. The latter is not completely impossible given that the magnetic field is found to inhibit the hole motion and to stabilize the Néel order even in nonSC $RBa_2Cu_3O_{6+x}$ crystals (Fig. 5). One can see that the Néel temperature indicated in the plot by an arrow, increases by $\sim 2$–$3$ K under 16.5 T field $H/|c|, H=\pm 16.5$ T. As a result, the corresponding resistivity upturn also starts from higher temperatures, which causes a significant magnetoresistance $\Delta \rho/\rho_c$ reaching $\approx 12\%$ at low $T$ (inset in Fig. 5). It is worth recalling that the resistivity measurements are incapable of providing information on $T_N$ values within the SC state and the data points in Fig. 3 located inside the SC region were obtained after the SC was killed by the 16.5 T magnetic field. If the field-induced increase in $T_N$ in the SC region is two-three times larger than in nonSC ones, the true magnetic state underlying the superconductivity in $R-123$ may well be a cluster spin glass which is driven to a long-range AF order only with the SC suppression.

B. AF order and the c-axis conductivity

After establishing the phase diagram, we can consider what is the “normal state” (at $T>T_N$, $T_c$) that both the AF and SC orders originate from and how it is affected by the AF ordering. At a first glance, the long-range antiferromagnetism just localizes the doped holes, thus strictly opposing the superconductivity (Figs. 1 and 5). However, there exists a narrow doping range on the phase diagram in Fig. 3 where the superconductivity sets in at $T<T_N$ (see also Figs. 1 and 2), which raises additional questions concerning the mechanism of hole transport in the AF state. In particular, does the c-axis transport actually acquire an activated character with establishing the long-range AF order? If the AF-ordered state is inevitably an insulator with an appreciable energy gap, how does the superconductivity emerge at $T_c<T_N$?

The simplest way of evaluating the activation energy is to plot the derivative $d \ln \rho_{\text{ao}} / d(1/T) = \Delta(T)$ which gives a $T$-independent $\Delta$ for ordinary activation behavior $\rho = A \exp(\Delta/T)$ or an effective gap $\Delta(T) \propto T^{-k}$ for variable range hopping $\rho = A \exp(B/T^{k})$. In contrast to expectations, the derivative data in Fig. 6 show a clear positive curvature ($k<0$) in the entire temperature region below $T_N$, which is inconsistent with a simple activated transport. The AF ordering therefore inhibits the hole motion, causing a sharp kink at $T_N$, but does not immediately turn the system into an insulator. With further reduction of the carrier density and increase of $T_N$ above 50–100 K, the hole motion indeed switches to the hopping regime, yet in the low-$T_N$ region shown in Figs. 5 and 6 the quasimetallic conduction survives, leaving some room for superconductivity. Strictly speaking, a possibility that the resistivity even in low-$T_N$ crystals in Figs. 5
and 6 will acquire an insulating behavior in the subKelvin region still remains, but the corresponding insulating gap in that case would be too small to affect the occurrence of superconductivity.

The derivative plots in Fig. 6 have other important implications. First, all the derivatives $d \ln \rho_c/d(1/T)$ at low $T$ appear to follow roughly the same curve until breaking down at $T_N$. The same low-$T$ resistivity behavior is observed even in superconducting crystals after suppressing SC with magnetic field (right panel). The $c$-axis transport mechanism in the spin-ordered state is therefore kept unchanged in the critical doping range of the AF-SC transformation; what changes with doping and magnetic field is merely the Néel temperature (Fig. 6). Second, the $\rho_c(T)$ upturn at low $T$—marked by the positive sign of the derivative in Fig. 6—is solely, or at least predominantly caused by the AF ordering. As the spin order melts above $T_N$, the resistivity quickly switches to the metallic behavior [negative sign of $d \ln \rho_c/d(1/T)$]. With decreasing $T_N$ the region of the resistivity upturn shrinks with a clear tendency to vanish completely for $T_N \to 0$ (Fig. 6). Therefore, in the $T_N=0$ limit (if it could be reached without getting into the SC state) the $c$-axis resistivity would keep the metal-like behavior down to low $T$ slowly decreasing and approaching a constant value at zero temperature.

The obtained data indicate that, in the absence of the interlayer AF order, the charge carriers in $\text{RBA}_2\text{Cu}_3\text{O}_{6+x}$ can move along the $c$ axis without thermal activation. Regardless of whether the nonactivated interlayer conduction occurs through a coherent, metallic transport, or by an energy-conserving tunneling, it requires itinerant quasiparticles to be already formed in CuO$_2$ planes. The latter in turn implies the in-plane charge transport in the nonAF-ordered state to be invariably metallic as well. Indeed, the in-plane resistivity $\rho_{ab}(T)$ was reported to exhibit a metallic behavior without any upturn until entering the Néel state. Consequently, the normal state playing a role of a starting point for the AF-SC phase competition in $\text{RBA}_2\text{Cu}_3\text{O}_{6+x}$ is an anisotropic 3D metal that experiences the impact of SC and AF ordering upon lowering the temperature.

**IV. DISCUSSION**

By now, many theoretical approaches have been suggested to describe the frustration that the doped holes introduce to the spin system and the resulting suppression of the long-range AF order. Namely, the doped holes were proposed to switch the exchange interaction between the nearest Cu spins from antiferromagnetic to ferromagnetic while being localized, or to continuously break the AF bonds while moving by the nearest-neighbor hopping. Spin distortions accompanying the doped holes are not restricted to the nearest Cu ions, but spread for several unit cells away and may form spin polarons, 2D spin vortices, spin spirals, AF domain walls, or complex spin textures controlled by the distribution of the dopants. Whatever the case, upon increasing the hole density the resulting spin distortions are suggested to gradually reduce the average staggered magnetization in each CuO$_2$ plane to zero, thereby making impossible the interlayer spin ordering and hence the long-range AF order.

While the mentioned above theoretical approaches have been positively tested on the prototype high-$T_c$ cuprate $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$, the considerably different behavior of $\text{RBA}_2\text{Cu}_3\text{O}_{6+x}$ may call for reconsidering the impact of doped holes on the long-range AF order. Indeed, the AF state in $\text{RBA}_2\text{Cu}_3\text{O}_{6+x}$ survives up to 2–3 times higher hole concentrations than in $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$, pointing to a significant material dependence of the AF-order suppression. Furthermore, the low-$T_N$ tail of the AF order spreading in the phase diagram toward SC compositions (Fig. 3) is hard to explain in the framework of those approaches, since they predict no slowing down of the spin frustration with increasing doping. One more difficulty for the existing models is the sharp AF transitions observed down to low $T_N$‘s, which are hard to reconcile with strongly frustrated spin states and disordered spin textures in CuO$_2$ planes suggested to emerge with doping.

Nevertheless, most of the observed features may be qualitatively accounted for by assuming the doped holes to break the long-range AF order in CuO$_2$ planes into 2D domains. Indeed, since the Néel temperature in that case should be determined by the interlayer spin coupling integrated over an average domain, a gradual reduction of the domain size with doping can explain why $T_N$ approaches zero value slowly and smoothly instead of vanishing abruptly. It is worth recalling that the interlayer coupling in $\text{RBA}_2\text{Cu}_3\text{O}_{6+x}$ is not frustrated in contrast to $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$ and thus should allow the Néel order in the former to survive up to much higher doping.

The idea of moving holes frustrating the AF order implies a reverse impact as well, which makes the origin of the observed antiferromagnetism-induced hole localization intuitively clear. As the penalty for frustrating the AF bonds exceeds the kinetic energy of a doped hole, the hole should get trapped by the environment. While in CuO$_2$ planes the hole hopping to the nearest-neighbor sites is inhibited smoothly due to the gradual development of AF correlations, the suppression of the $c$-axis hole hopping should occur abruptly with the $c$-axis spin ordering at $T_N$ exactly as observed experimentally. However a question remains: when exactly does the charge transport switch its character from metallic to nonmetallic, i.e., thermally activated hopping? A comparison of the $c$-axis conductivity for the cases of spins being or not AF ordered along that axis has demonstrated that in the absence of spin order the holes move without thermal activation. Therefore, the normal state in the AF-SC transformation region (at temperatures above $T_N$ and $T_c$) is certainly a 3D metal. As the AF order develops, the holes are driven to localization, yet the exact position of the metal-insulator transition (MIT) remains unsettled.

Previous attempts to determine the MIT location on the phase diagram of $\text{RBA}_2\text{Cu}_3\text{O}_{6+x}$ employed both the charge-transport and heat-transport measurements, yet resulted in somewhat controversial conclusions; namely, the MIT was suggested to occur at hole densities lower, equal, or even higher than the SC onset (see Refs. 22 and 23 and references therein). Is there a universal doping level for MIT in $\text{RBA}_2\text{Cu}_3\text{O}_{6+x}$ at all? An important result following from the present study is that the metal-insulator transformation is largely driven by the AF spin ordering. Given that the AF
correlations in $RBa_2Cu_3O_{6+x}$ are not determined solely by the hole density, but develop with decreasing temperature, get stabilized by magnetic fields, and presumably depend on some material parameters, there is little reason to expect the MIT to be located on the phase diagram at some universal hole density. Instead, the metal-insulator transformation may take place depending on doping, magnetic field, material parameters, and, in a sense, even on temperature. The latter means that the hole localization gains strength on cooling due to the development of the AF order and thus the metal-insulator border may be crossed at some transition temperature. What still remains to be understood is how the magnetic field enters the AF-SC interplay. One possibility is that the magnetic field affects both orders almost independently, suppressing superconductivity, and stabilizing the long-range AF state. Indeed, a magnetic field applied to a 2D Heisenberg antiferromagnet can introduce anisotropy to the spin system, thus suppressing the spin fluctuations and allowing the long-range Néel order to be established. Alternatively, the AF and SC ground states may be separated by a first-order transition, as were predicted, for instance, in the SO(5) theory and the magnetic field may directly affect their competition. Microscopic probes capable of detecting the AF order inside the superconducting state are indispensable for this issue to be clarified.

**ACKNOWLEDGMENTS**

We thank V. F. Gantmakher and D. V. Shovkun for fruitful discussions and acknowledge support by RFBR (Grants No. 05-02-16973 and No. 09-02-01224) and the integration project SB RAS No.81.

12. A pronounced crossover was reproducibly observed in the $c$-axis resistivity of underdoped $RBa_2Cu_3O_{6+x}$ crystals (Refs. 8–11) and some other layered oxides (Ref. 26), yet its origin is still insufficiently understood.