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# Pressure-Induced Transformations in Two-Dimensional Polymeric Phases of $C_{60}$ <sup>1</sup>

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**Abstract**—The structural stability of the tetragonal and rhombohedral two-dimensional (2D) polymeric phases of  $C_{60}$  was studied under pressures up to 27 GPa at room temperature by means of in situ Raman scattering spectroscopy. The results show that the tetragonal 2D phase undergoes an irreversible transformation in the region of 20 GPa while no pressure-induced transitions were observed for the rhombohedral 2D phase. The obtained data are discussed within the framework of recent numerical calculations, which predict the pressure-induced transformation of the 2D polymeric phases of  $C_{60}$  into three-dimensional (3D) polymers in the pressure range 14–20 GPa. © 2002 MAIK “Nauka/Interperiodica”.

The pristine  $C_{60}$  in the condensed phase has great potential for polymerization due to the existence of 30 double C=C bonds in the fullerene molecular cage. Covalent polymeric bonds between adjacent molecules are usually built up by the [2 + 2] cyclo-addition mechanism, which results in an  $sp^3$ -like four-fold coordination of a number of carbon atoms of the molecular cage [1]. The crystal structures of the known polymeric phases of  $C_{60}$  were identified as one-dimensional (1D) orthorhombic, 2D tetragonal, 2D rhombohedral, and 3D cubic [2–4]. These phases, as well as the so-called hard fullerite phases of 3D polymeric nature [5], can be prepared under various conditions of high pressure and high temperature treatment. Recent theoretical calculations have predicted that the 2D polymeric phase of  $C_{60}$  can be transformed, at room temperature, by uniaxial compression into a 3D polymeric phase [6]. This polymerization will take place at  $P \approx 20$  GPa and results in the formation of a stable metallic phase having 24  $sp^3$ - and 36  $sp^2$ -like coordinated carbon atoms in each  $C_{60}$  molecule. Another theoretical study [7] predicts that uniaxial compression up to  $\sim 14$  GPa perpendicular to the chains in 1D or the planes in 2D polymers of  $C_{60}$  leads to 3D polymerization with 52, 56, or even 60  $sp^3$ -like coordinated carbon atoms per molecule. To explore these predictions we studied the pressure response of the 2D tetragonal and rhombohedral polymeric phases of  $C_{60}$  at pressures up to  $\sim 27$  GPa by means of in situ Raman scattering using the diamond anvil cell (DAC) technique.

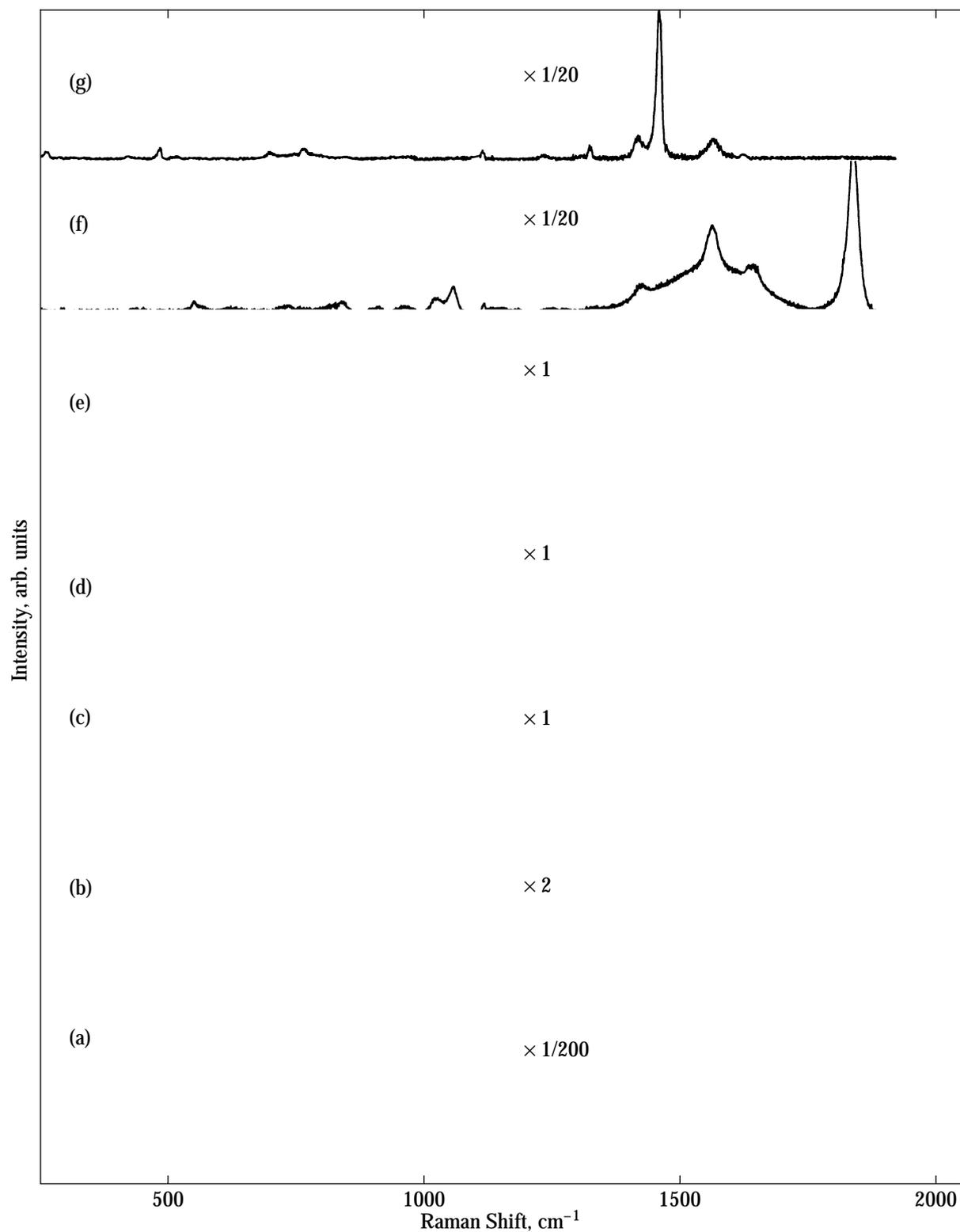
<sup>1</sup> This article was submitted by authors in English.

## 1. EXPERIMENTAL

The 2D polymeric phases of  $C_{60}$  were obtained by high pressure and high temperature treatment of 99.99% pure pristine  $C_{60}$  at  $T = 820$  K  $P \approx 2.5$  GPa and  $T = 773$  K,  $P \approx 5$  GPa for the tetragonal and rhombohedral polymers, respectively. Raman spectra were recorded using a triple monochromator (DILOR XY-500) equipped with a CCD liquid-nitrogen-cooled detector system. The 514.5 nm line of an Ar<sup>+</sup> laser with a beam power of 2–10 mW, measured before the cell, was used for excitation. Measurements of the Raman spectra at high pressure were carried out using a DAC of the Mao–Bell type. A 4 : 1 methanol-ethanol mixture was used as the pressure transmitting medium, and the ruby fluorescence technique was used for pressure calibration. The specimens had dimensions of  $\sim 100$   $\mu$ m and were selected from a batch material by checking their Raman spectrum such that they were spatially uniform and typical of the corresponding phase.

## 2. RESULTS AND DISCUSSION

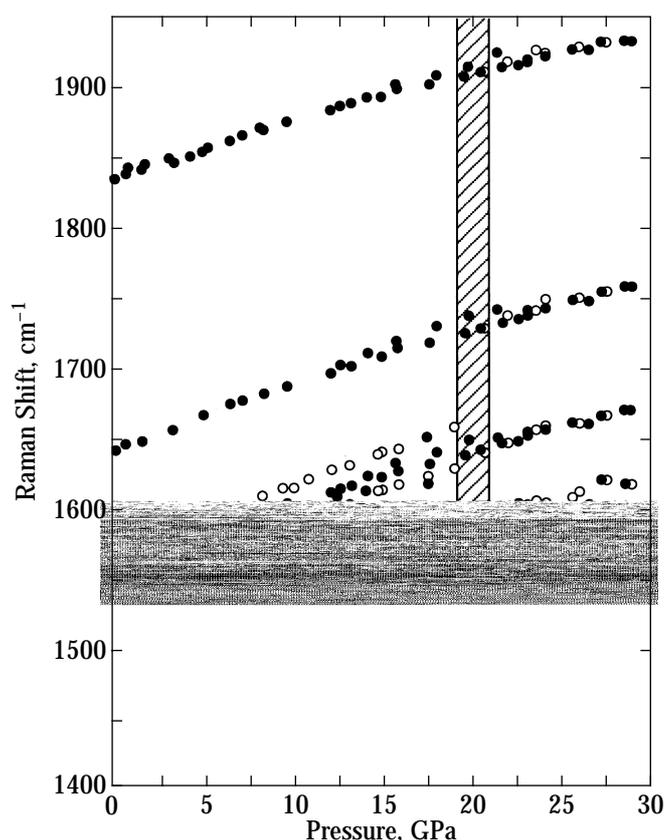
The Raman spectra of the 2D tetragonal polymer of  $C_{60}$  pressures up to  $\sim 27$  GPa and room temperature are illustrated in Fig. 1. The initial spectrum (Fig. 1a), taken under normal conditions, is identical to the spectra of this material reported earlier [8]. The spectrum is richer than that of the pristine  $C_{60}$  due to the lower symmetry of the molecule, which, in turn, results in the splitting of degenerate phonon modes and/or the activation of silent modes. As the pressure increases, the Raman peaks shift to higher energies and their bandwidth gradually increases (Fig. 1b,  $P \approx 14.1$  GPa). Dra-



**Fig. 1.** Raman spectra of the tetragonal 2D polymeric phase of  $C_{60}$  at high pressure and room temperature. Spectra (a), (b), (c), (d) are related to the increase in pressure from ambient to 14.1, 20.7, and 27.5 GPa, respectively. Spectra (e) and (f) are related to the decrease in pressure from 12 GPa to ambient. Spectrum (g) is related to the principal part of the detonated sample after the pressure is released.

matic changes are first recorded at a pressure of  $\sim 20.7$  GPa (Fig. 1c). New features appear which grow rapidly in intensity with increasing pressure up to the highest pressure attained in this study (Fig. 1d,  $P \approx 27.5$  GPa). The transformation of the Raman spectrum is associated with the appearance of new intense peaks near  $\sim 1000$  and  $\sim 1920$   $\text{cm}^{-1}$ , as well as the disappearance of some old ones. The total number of Raman active modes in the new spectrum is smaller than, whereas the intensity and position of other peaks differ considerably from, those in the initial spectrum. The majority of the peaks can be tracked back to the peaks observed in the tetragonal polymeric phase of  $\text{C}_{60}$  and may be related to the fullerene molecular cage. The decrease in pressure (Fig. 1e,  $P \approx 12$  GPa) up to its total release to ambient pressure (Fig. 1f) results in a gradual shift of the Raman peaks to lower energies, as well as in preservation of the spectrum to ambient pressure. The pressure dependence of the Raman modes of the tetragonal 2D polymeric phase of  $\text{C}_{60}$  in the high-energy region is shown in Fig. 2. The shaded area marks the pressure region around  $\sim 20$  GPa where drastic changes in the Raman spectrum took place. Similar changes, at  $\sim 20$  GPa, are shown also in the pressure dependence of almost all Raman modes. The decrease in pressure from  $\sim 27$  GPa to ambient pressure, in a time period of about two weeks, does not lead to any observable changes in the Raman intensity distribution and the new high-pressure phase remains stable. It is important to note that the recovered sample was in fact in a metastable phase, which was detonated by the probing laser beam during the Raman measurements after a time period of a few days from the moment of its exposure to air. The main part of the detonated sample is related to pristine or partially dimerized  $\text{C}_{60}$  (Fig. 1g), but there was also a small part related to "diamond-like graphite" that has strong Raman peaks at 1339 and 1595  $\text{cm}^{-1}$  resembling those of diamond and graphite, respectively.

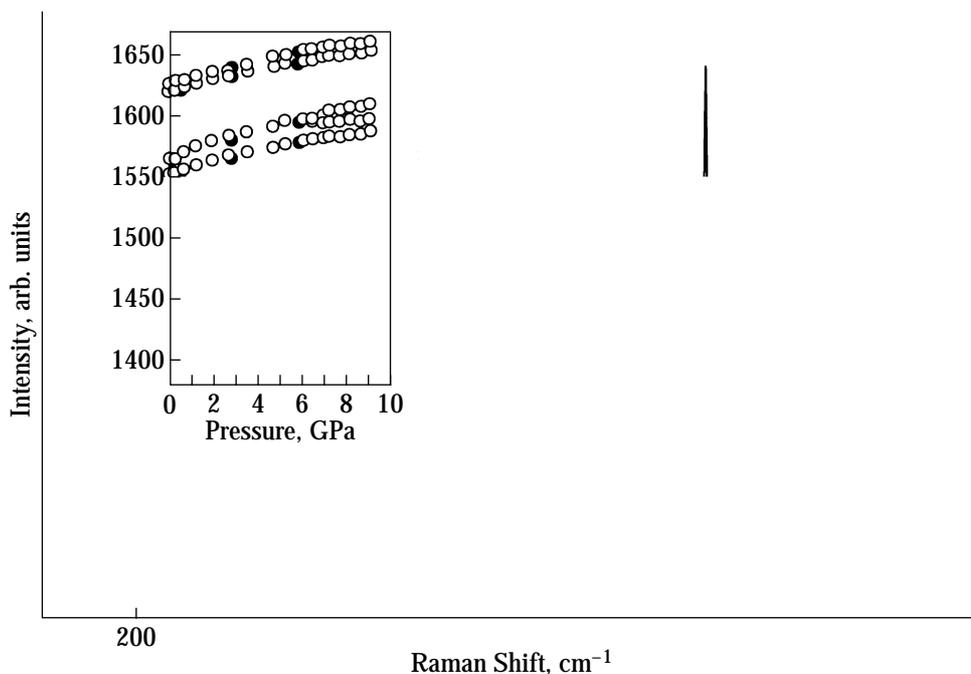
The Raman spectrum of the rhombohedral 2D polymeric phase of  $\text{C}_{60}$ , under normal conditions, is shown in Fig. 3. As in the case of the other polymeric phases of  $\text{C}_{60}$ , the lowering in symmetry results in a very rich Raman spectrum, especially in the low energy region [9]. The spectrum differs in many details from that of the tetragonal 2D polymer of  $\text{C}_{60}$ . An important difference is related to the frequency of the pentagonal-pinch (PP) mode, which is gradually softened in the polymeric phases of  $\text{C}_{60}$  in accordance with the increasing number of the  $sp^3$ -like coordinated carbon atoms per  $\text{C}_{60}$  molecule. Thus the frequency of the PP-mode decreases from 1469  $\text{cm}^{-1}$  for pristine  $\text{C}_{60}$  to 1465, 1450, and 1408  $\text{cm}^{-1}$  for dimerized  $\text{C}_{60}$ , tetragonal, and rhombohedral polymeric phases, respectively. The pressure dependence of the Raman modes of the rhombohedral polymeric phase of  $\text{C}_{60}$ , studied up to 10 GPa, is shown in the insert to Fig. 3. The behavior of the



**Fig. 2.** Pressure dependence of the Raman frequencies of the 2D tetragonal polymeric phase of  $\text{C}_{60}$  in the high-energy region. Open (closed) symbols are related to the increase (decrease) in pressure. The shaded area near 20 GPa indicates the pressure region of the irreversible transformation of the material.

material is reversible with pressure (open and closed symbols indicate increases and decreases in pressure, respectively) and does not show any peculiarity in the pressure range investigated.

The obtained experimental data provide a strong indication that the tetragonal 2D polymeric phase of  $\text{C}_{60}$  undergoes a phase transition at  $\sim 20$  GPa. The transformation is irreversible and takes place via an intermediate disordered pretransitional state. The retention of the phonon modes of the  $\text{C}_{60}$  cage and the reduction of the total number of Raman active modes in the high-pressure phase are indications that this phase may be related to a 3D polymerized phase. This agrees well with the expected higher symmetry for the 3D polymerized phase. Another argument in favor of the 3D polymerization is the metastability of the high-pressure phase, which is transformed after a period of time into pristine or partially dimerized  $\text{C}_{60}$ . The most intriguing feature of the new phase in its Raman spectrum is the presence of a very intense peak at  $\sim 840$   $\text{cm}^{-1}$ , which cannot be related to any of the fullerene molecular cage modes. The appearance of this peak may be associated



with the considerable distortion of the PP-mode of the pristine  $C_{60}$  molecule due to the breakdown of a large number of double C=C bonds and the  $sp^3$ -like hybridization of the carbon atoms involved in the intermolecular covalent bonding. The similar string Raman peak near  $1840\text{ cm}^{-1}$ , in some chemical compounds of carbon, is related to the stretching vibration of the isolated double  $>C=C<$  bond [10]. The existence of the  $1840\text{-cm}^{-1}$  mode in the high-pressure phase is note compatible with the theoretical calculations in [7], as they predict no phonon mode above  $1600\text{ cm}^{-1}$ . On the other hand, the observed transition pressure near  $\sim 20$  GPa agrees well with the theoretical predictions in [6]; therefore, our results point to the polymerization mechanism predicted in that study.

In conclusion, the pressure-induced irreversible transformation of the 2D tetragonal polymeric phase of  $C_{60}$  was observed at  $\sim 20$  GPa by means of in situ Raman scattering study. The data obtained provide a strong indication that the new phase is related to a 3D polymeric phase of  $C_{60}$ . The Raman study of the rhombohedral polymeric phase of  $C_{60}$  at pressures up to 10 GPa does not show any structural phase transition or irreversible transformation of the material.

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