Abstract: High-pressure torsion has been used to obtain the ultra-fine grained (UFG) state with a high specific area of grain boundaries (GBs) in Al-Zn, Al-Mg, Cu-Ag, Cu-Co, and Cu-Ni solid solutions with face-centered cubic (fcc) lattices. The UFG samples were heated in a differential scanning calorimeter (DSC). Small endothermic peaks in the DSC curves were observed in the one-phase solid-solution area of the respective phase diagrams, i.e., far away from the bulk solidus and solvus lines. A possible explanation of these endothermic peaks is based on the hypothesis of phase transformations between GB complexions. This hypothesis has been supported by observations with transmission electron microscopy and electron backscattering diffraction. The new lines of GB phase transformations have been constructed in the Al-Zn, Al-Mg, Cu-Ag, Cu-Co, and Cu-Ni bulk phase diagrams.

Keywords: grain boundaries; high-pressure torsion; differential scanning calorimeter; phase transformation; metastable phases; Al-Zn; Al-Mg; Cu-Ag; Cu-Co; Cu-Ni

1. Introduction

Phase transformations can take place not only in the volume of a material (as 3D transformations) but also on free surfaces or in the interphase boundaries (IBs) and grain boundaries (GBs). Such unusual phase transformations in GBs were theoretically discovered for the first time by Cahn [1] and later by Ebner and Saam [2]. Different reviews discuss the experimental works done after the predictions of References [1,2], and various aspects of GB phase transitions have been studied [3–13]. As one of the important results of these works, the new additional lines of GB phase transformations appeared in conventional 3D-phase diagrams [14–27]. It was quite simple to study the GB phase transformations in the two-phase (or multiphase) regions of the bulk phase diagrams. Important examples of such phase transformations are the GB wetting–dewetting ones [8,28–32]. The thickness of a layer of a GB wetting phase is not small, about a few micrometers. In other words, the second phase wetting for a GB is nothing other than a “normal” bulk phase. However, the GB itself disappears after transition from partial to complete wetting. Therefore, the transition from incomplete to complete GB wetting is indeed a GB transition because a GB ceases to exist after the transition. The GB wetting transformations are easy...
to observe even when using optical microscopy [33–35]. The GB phase transitions were also observed in the one-phase regions of the conventional 3D-phase diagrams. Thus, the GB premelting or prewetting can take place in the bulk solid solution [8,18,36–45]. In the case of premelting or prewetting, the GBs can contain a few-nanometers-thin layer of a second phase which can be stable in GB but cannot be thermodynamically stable in bulk [18,46–48]. Such intergranular films (or IGFs) are not easy to observe, mainly because they are very thin (a few nanometers). Therefore, experimental methods for the investigation of such IGFs only became available recently. These include, for example, high-resolution transmission electron microscopy (HRTEM) in the aberration-corrected microscopes [17,47,49–52] as well as three-dimensional atom probe tomography (APT) [10,48,53,54].

Quite frequently the GB phase transitions occur above room temperature (RT). As a result, one needs to quench the sample to study it further with the aid of HRTEM or 3-D APT. However, one cannot be sure that the GB structure after quenching is the same as before quenching. It is more complicated to quench the high-temperature structure of metallic alloys than that of nitrides or oxides [9,23,55–63]. Therefore, it is easy to understand why the thin IGFs were first observed in silicon nitride [64]. The quenching of metallic alloys seldom allows the “freezing” of high-temperature IGFs or their transformed remnants [10,22,24,47,48,65–68]. In the meantime, different kinds of GB phases or IGFs have been discovered together with respective GB transformations [18–20,62,66,67,69–72]. According to the work of S.J. Dillon et al. [73] GB phases or IGFs are called GB complexions [11,17,18,21,22,44,47,62,63,67,72–79]. We have to underline here the subtle differences between the terms “GB phases”, “IGFs”, “GB complexions” and respective GB transitions, on the one hand, and the GB transformations with the participation of the bulk phase, such as GB wetting transformations, on the other hand.

Advantageously, if one cannot quench the GB complexions to RT to study them ex-situ, one can investigate GB phase transformations indirectly. This can be performed because the GB complexions and GB transformations manifest themselves in changes in GB properties. In turn, the GB properties influence the properties of the polycrystal as a whole [54,80,81]. These changes in properties make it possible to exploit in-situ studies. In other words, for the measurements at actual temperatures of GB transitions. For example, one can measure the GB migration rate [23,82–84]. The discontinuities in the Arrhenius plots of GB mobilities mark the GB phase transitions. Usually, a GB containing a liquid-like IGF becomes more mobile in comparison with a “dry” one, such as, for example, Al/Al GBs doped with Ga [84]. Similarly, changes in GB diffusivity can be observed [85–88]. For example, the tracer self-diffusion of copper and bismuth are drastically (by about two orders of magnitude) accelerated when Cu/Cu GBs obtain the liquid-like Bi-rich premelting layer [88]. Another important example is the chemical zinc diffusion along individual GBs in Fe-Si bicrystals with various silicon contents [86,87]. The GB diffusivity of zinc is accelerated by about two orders of magnitude when a GB in doped iron becomes the premelting layer. GB faceting/roughening and GB energy [60,78,89–94], as well as GB sliding [95–99], are also very sensitive to the formation of thin IGFs. For example, at the point of GB premelting transition in a Cu-Bi system, the first derivative of GB energy has a break [91]. This, thus, directly demonstrates that the GB premelting transition is of the first order. The GB faceting in a Cu-Bi system also follows this change [93].

Thus, due to the GB phase transformations, the GBs can become very ductile [100–102] or very brittle [103–105]. The ductility increases if the GB phase is plastic and serves as a kind of lubricant for the abutting grains [100–102]. If the GB layers are brittle as in Ni-Bi alloys, superalloys or molybdenum alloys [103–105], the overall brittleness of the polycrystal increases as well. GB phase transformations can also lead to a change in GB segregation [49–51,53,66,70,74,104–108]. The most prominent and investigated examples are the (already mentioned) Cu-Bi and Ni-Bi alloys or oxides, such as ytterbium-doped magnesium aluminate spinel.

One can also investigate the GB-controlled properties of polycrystals, such as creep [100], or superductility [109–119]. The UFG alloy becomes superductile if the thin film of a liquid-like phase lubricates the GB, for example, in ternary Al-Zn-Mg alloys close to the bulk liquidus lines [115–119].
However, the superductility also appears if the GB phase is solid [109–114]. It only needs to be ductile enough to lubricate the GB between abutting matrix grains (like zinc-rich layers between aluminum grains) [109–114]. In any case, the mechanical properties of the GB phase strongly influence the overall strength of a material [76,120–122]. Since GB layers influence the mobility of individual GBs, they also change the migrational ability of the GB ensemble in a polycrystal and, thus, the grain growth [54,56–59,75,76,79,80,83,106,123–130]. Most pronounced here is the transition from normal to abnormal grain growth (i.e., from single modal to the bi- or multimodal distribution of grain size). Similarly, the presence of liquid-like GB layers can increase the diffusion permeability of individual GBs. If the portion of such GBs in a polycrystal is high enough (and the percolation threshold is exceeded), the rate of densification during sintering can significantly increase [26,57,81,131–136]. If the GBs possess low electric conductivity and their portion is high (again, the percolation threshold is exceeded) the overall conductivity of the polycrystal decreases [9,95,137–141]. On the other hand, the polycrystal can become ferromagnetic (for example, nanograin zinc oxide) if the amount of “magnetic” GBs is high enough [142–144]. The graphene-modified GB complexion in skutterudite can improve the thermal conductivity [145] similar to GB layering transitions [69].

The phase transitions between bulk phases can be endo- or exothermal. As a result, they can be observed and investigated by thermal methods, such as differential scanning calorimetry (DSC) [146,147]. In fine-grained (or better in ultrafine-grained (UFG)) polycrystals, the specific area of GBs can be very high. In addition, the specific volume of GB films in such materials is not low when compared with the total volume. As a result, one can measure the thermal effect of GB phase transitions with the aid of a DSC [114,148–150]. DSC also permits to observe the indications of GB phase transformations in nanocrystalline ceramics [151–160]. The GB phase transformations can also lead to abnormal grain growth and radical changes in the GB character distribution [54,56–59,75,76,79,80,83,106,123–130,151,154–159].

Therefore, the goal of this work is to study GB phase transformations and GB complexions (i) in-situ; (ii) in poorly quenchable metallic alloys; (iii) deep in the solid-solutions area (i.e., removed from the solidus or solvus lines); (iv) by comparing the samples with low and high specific area of GBs; and (v) by the use of thermal effect of such transformations and abnormal grain growth. The choice of alloys (Al-Zn, Al-Mg, Cu-Co, Cu-Ag, and Cu-Ni) is primarily related to the wide technological application of these materials, such as aircraft construction, mechanical engineering, construction, as well as for solid oxide fuel cells interconnect material.

2. Experimental

Aluminum alloys (Al-20 wt.% Zn, Al-10 wt.% Mg) and copper alloys (Cu-2.2 wt.% Co, Cu-4.9 wt.% Co, Cu-77 wt.% Ni, Cu-2.8 wt.% Ag, Cu-4 wt.% Ag, Cu-5 wt.% Ag, Cu-6.5 wt.% Ag) were manufactured of highly pure components (5N5 Al, 5N Zn, 4N5 Mg, 5N Cu, 5N Ag, 4N5 Ni, and 5N Co) using vacuum induction melting (Institute of Solid State Physics, Chernogolovka, Russia) in the form cylindrical ingots with a diameter of 10 mm. For further research, disks with a thickness of 3.5 and 0.7 mm were cut from the ingots by spark erosion. The 0.7 mm thick disks were subjected to the high-pressure torsion (HPT) after chemical etching using a custom-built computer controlled HPT device (W. Klement GmbH, Lang, Austria) at room temperature at a pressure of 5 GPa (5 torsions, 1 rpm). After HPT, the grain size in the samples became very fine (see Figure 1). All samples for structural and calorimetric investigations were cut from the deformed disks at a distance of 3 mm from the sample centre in the HPT-treated samples. The samples were heated in the differential scanning calorimeter (DSC 404 F1 Pegasus, Netzsch and Du Pont 910, TA Instruments (New Castle, UK) from room temperature to melting point at a rate of 10 K/min. After heating the samples were cooled in the calorimeter to room temperature at the rate of about 10 K/s. After HPT, the Cu-77 wt.% Ni samples were sealed in the evacuated quartz ampoules (residual pressure of 10^-4 Pa) and annealed at 600 °C, 3 h; 750 °C, 2 h, and 900 °C, 1 h. The grain structure of the annealed Cu-77 wt.% Ni samples was studied with the aid of the electron backscattering diffraction (EBSD) (FEI, Hillsboro, OR, USA) technique.
EBSD maps and misorientation angle charts were obtained using FEI QUANTA 3D FEG scanning microscope equipped with EBSD TSL hardware at the acceleration voltage of 20 kV and tilt angle of 70°. Post-processing of the acquired data was performed by the use of TSL OIM Analysis 7.30 software. Transmission electron microscopy (TEM) (TECNAI, Hillsboro, OR, USA) investigations were conducted using a TECNAI G2 F20 microscope at an accelerating voltage of 200 kV. Prior inspection of the obtained material was carried out also using an FEI E-SEM XL30 scanning electron microscope equipped with EDAX Genesis energy-dispersive X-ray spectrometer (EDS) (FEI, Hillsboro, OR, USA). X-ray diffraction (Philips, Manchester, UK) (the XRD patterns were obtained using Bragg-Brentano geometry in a powder diffractometer Philips X’Pert with Cu-Kα radiation). The detailed microstructure studies were performed by the use of TECNAI G2 FEG super TWIN (200 kV) transmission electron microscope (TEM) equipped with an energy dispersive X-ray (EDS) spectrometer manufactured by EDAX (FEI, Hillsboro, OR, USA). The thin foils for TEM observation were prepared by a twin-jet polishing technique using an electrolyte D2 manufactured by Struers company.

![Figure 1](image1.png)

**Figure 1.** Transmission electron microscope (TEM) micrographs of samples after high-pressure torsion (HPT)-treatment and before heating in differential scanning calorimeter (DSC) calorimeter. (a) Cu-5 wt.% Co, bright field; (b) Cu-2.8 wt.% Ag, dark field, (c) Cu-77 wt.% Ni, dark field. Some microtwins are shown by the arrows.

### 3. Results

Figure 2 shows the bulk phase diagrams of the eutectic systems Al-Zn (Figure 2a), Al-Mg (Figure 2b), Cu-Ag (Figure 2c), peritectic system Cu-Co (Figure 2d), and system Cu-Ni with continuous raw solid solutions and a decomposition dome at low temperature (Figure 2d) [161]. Vertically arranged experimental points correspond to the compositions Al-20 wt.% Zn (Figure 2a), Al-10 wt.% Mg (Figure 2b), Cu-2.2 wt.% Co, Cu-4.9 wt.% Co (Figure 2c), Cu-77 wt.% Ni (Figure 2d), Cu-2.8 wt.% Ag, Cu-4 wt.% Ag, Cu-5 wt.% Ag, Cu-6.5 wt.% Ag (Figure 2e).

![Figure 2](image2.png)

**Figure 2. Cont.**
In Figure 3 the temperature dependences of the heat flow (DSC curves) are shown for the Al-20 wt.% Zn, Al-10 wt.% Mg, and Cu-4.9 wt.% Co alloys before and after HPT-treatment. After HPT, the grain size in Al-based alloys was about 500 nm [109,113,162], and in Cu-based alloys, it was 150–300 nm [163]. Before HPT, the Al-based alloys contained the supersaturated (Al) solid solution and, respectively, the particles of Zn or Al$_2$Mg$_3$ intermetallic phase [109,113,162]. The Cu-Co alloys contained supersaturated (Cu) solid solution and particles of Co [163] before HPT. After HPT, the composition of a solid solution and amount of the second phase changed slightly, but the mean effect important for this work was the grain refinement. The heating of these ultra-fine grained (UFG) HPT-treatment samples in a DSC apparatus led to the growth of grains up to about 30–50 μm. The grain size in the coarse-grained (CG) as-cast samples studied for comparison was between 0.1 and 5 mm. Thus, we compared the heat flow in samples with a low and high specific area of grain boundaries.

DSC curves in Figure 3 have a complex shape with overlapping peaks. We used the standard procedure for the quantification of the DSC curves [164] (these procedures are also included in the quantification of the software of modern DSC equipment). According to this approach, the position of the deep minimum (peak) corresponds to the melting temperature, i.e., the liquidus temperature in the bulk phase diagram. The position of this deep minimum (peak) can be slightly influenced by the heating rate. Therefore, we used the same heating rate in all experiments. The onset temperature of the peak corresponds to the point of intersection of the tangent drawn through the inflexion point of the curve with the extrapolated baseline. The baseline is a virtual line drawn through the interval in which the phase transition occurs.

**Figure 2.** The bulk phase diagrams of the systems (a) Al-Zn; (b) Al-Mg; (c) Cu-Co; (d) Cu-Ag; and (e) Cu-Ni with lines of bulk phase transitions [151]. $T_w$—grain boundaries (GB) wetting phase transformation by liquid phase, $T_{ws}$—GB wetting phase transformation by solid phase.
The temperature of these endothermic peaks is shown by the crosses in Figure 2a. We can observe that they are “deep” in the single-phase solid-solution region (Al) of the Al-Zn the bulk phase diagram. These peaks are more pronounced in the samples after HPT-treatment, i.e., in samples with smaller grain size containing, therefore, more grain boundaries than CG samples. The temperature of these deep minimum (peak) can be slightly influenced by the grain size in the bulk phase diagram. The position of this deep minimum (peak) corresponds to the melting temperature, i.e., the liquidus temperature.

For the Al-20 wt.% Zn alloy (Figure 3a) small endothermic peaks are observed in the DSC curves between 350 to 500 °C. These peaks are more pronounced in the samples after HPT-treatment, i.e., in samples with smaller grain size containing, therefore, more grain boundaries than CG samples. The temperature of these endothermic peaks is shown by the crosses in Figure 2a. We can observe that they are “deep” in the single-phase solid-solution region (Al) of the Al-Zn the bulk phase diagram (Figure 2a). Figure 4a,b shows microstructures of the Al-20 wt.% Zn alloy after annealing in DSC calorimeter and subsequent rapid cooling. The CG sample was annealed at $T = 350$ °C and UFG
sample (obtained by the HPT-treatment) at \( T = 320 \, ^\circ\text{C} \). In both samples, the precipitates of the second phase with a zinc content up to 33 wt.% Zn are visible in the GBs.

![Figure 4. TEM micrographs: (a) Al-20 wt.% Zn alloy before HPT-treatment at \( T = 350 \, ^\circ\text{C} \); (b) Al-20 wt.% Zn alloy after HPT-treatment at \( T = 320 \, ^\circ\text{C} \); (c) Al-10 wt.% Mg before HPT-treatment at \( T = 400 \, ^\circ\text{C} \); (d) Al-10 wt.% Mg alloy after HPT-treatment at \( T = 400 \, ^\circ\text{C} \).](image)

In the CG and UFG samples of the Al-10 wt. % Mg alloy, we also observed several small endothermic peaks in the DSC curves (Figure 3b) above 400 °C. These endothermic peaks are also more pronounced in UFG samples (obtained by HPT) when compared with their CG counterparts. The position of these peaks is shown by crosses in the Al-Mg phase diagram (Figure 2b). Similar to the Al-Zn system, these points are “deep” in the single-phase (Al) region of the phase diagram. In Figure 4c,d, the TEM micrographs for CG and UFG samples of Al-10 wt.% Mg alloy are shown. These CG and UFG samples were annealed in DSC spectrometer at \( T = 400 \, ^\circ\text{C} \) (i.e., close to that of the DSC peaks) and quickly cooled. GBs both in the CG and UFG samples contain fine precipitates of \( \text{Al}_3\text{Mg}_2 \) phase [31,148].

The DSC curves for CG and UFG Cu-Co alloys (Figure 3c) contain a pronounced additional peak for all samples at around \( T = 1070 \, ^\circ\text{C} \) (also denoted by crosses in the single-phase region of the Cu-Co phase diagram, Figure 2c). Figure 3c shows the DSC curves for the CG and UFG samples of Cu-4.9 wt.% Co alloy. The DSC curve of CG sample was decomposed into three endothermic peaks (at temperatures of 1069.8 °C, 1091.1 °C, and 1101 °C), and the DSC curve for UFG sample was decomposed into four peaks (at temperatures 1068.4 °C, 1088.2 °C, 1095.9 °C, and 1108.2 °C). The peaks in the DSC curves of CG sample (at \( T = 1091.1 \, ^\circ\text{C} \) and 1101 °C) and UFG sample (at \( T = 1095.9 \, ^\circ\text{C} \) and 1108.2 °C) correspond to the phase transformations occurring in the bulk in the two-phase region L + (Cu). The second peak (at \( T = 1088.2 \, ^\circ\text{C} \)) on the DSC curve of UFG sample lies below the line of the bulk solidus and corresponds to the melting of GBs. We assume that the first peaks (at \( T = 1069.8 \, ^\circ\text{C} \) and 1068.4 °C) in
both DSC curves correspond to the precipitates of the second phase (Co) as evidenced indirectly by
the data obtained in Reference [164].

The DSC curves for UFG Cu-2.8 wt.% Ag, Cu-4 wt.% Ag, Cu-5 wt.% Ag and Cu-6.5 wt.% Ag
alloys are shown in Figure 3d. These UFG samples were heated to complete melting, then cooled and
heated again to complete melting. The DSC curves for the second heating are shown in Figure 3e.
Differently from the UFG samples in the first heating, they were coarse-grained and contained
fewer GBs. The endothermal melting peak in all UFG samples had a complicated form and can
be interpreted as consisting of two or even three overlapping peaks. The respective tangents are shown
by thin lines. The DSC curves for the second heating, i.e., for coarse-grained samples, are more simple
when compared with DSC curves for UFG samples. We can see that only the curves for Cu-5 wt.% Ag
and Cu-6.5 wt.% Ag alloys have complicated melting peaks and the curves for the other two alloys
(with less silver) do not. The tangents in the main melting peak correspond to the liquidus and solidus
temperature and coincide well with each other for the first and second heating with only a couple
of degrees difference. The respective points are shown in the Cu-Ag phase diagram (Figure 2e) with
up- and down- full triangles. The additional melting peaks are more pronounced in the UFG samples
and lie at about 10–30 ºC below the bulk solidus line. They are shown in the Cu-Ag phase diagram
(Figure 2e) by open circles and squares. The difference between bulk and GB solidus increases with
increasing silver content from ~10 ºC to ~30 ºC (Figures 3d and 2e).

The DSC curves for UFG and CG Cu-5 wt.% Ag and Cu-6.5 wt.% Ag alloys contain a small
endothermic peak at around \( T = 778 \) ºC (Figure 3d,e). The DSC curve for the Cu-4 wt.% Ag alloy
contains such a peak in the second heating (Figure 3e). This peak is not visible in both the DSC curves
for UFG Cu-2.8 wt.% Ag alloy. These peaks are denoted by crosses in the single-phase region of the
Cu-Ag phase diagram (Figures 2e and 3d). The thermal effect of this peak increases with increasing
silver content.

The DSC curves for UFG Cu-2.8 wt.% Ag, Cu-4 wt.% Ag, Cu-5 wt.% Ag and Cu-6.5 wt.% Ag alloys contain a small endothermic peak at around \( T = 778 \) ºC (Figure 3d,e). The DSC curve for the Cu-4 wt.% Ag alloy contains such a peak in the second heating (Figure 3e). This peak is not visible in both the DSC curves for UFG Cu-2.8 wt.% Ag alloy. These peaks are denoted by crosses in the single-phase region of the Cu-Ag phase diagram (Figures 2e and 3d). The thermal effect of this peak increases with increasing silver content.

The UFG samples of Cu-77 wt.% Ni alloy were annealed at 600 ºC, 3 h; 750 ºC, 2 h, and 900 ºC, 1 h).
Unique grain color maps acquired by EBSD technique (Figure 5a,c,e) show the resultant grain structure
after extensive grain growth. Micrographs in Figure 5a,c permit to observe that the distribution of small and large grains is not uniform. It indicates abnormal grain growth in the sample. The mean grain size is 1.6, 6.8, and 18.3 µm respectively. Numerous flat twin (or \( \Sigma 3 \)) GBs are clearly visible on the maps. Figure 5b,d,f show the misorientation angle charts for these polycrystals based on unique grain color maps. These polycrystals have about 0.3–0.7% of GBs with low misorientation angle (i.e., below 5º). The number of GBs with a misorientation angle between 5º and 58º fluctuates from 72% at 600 ºC, 3 h through 55% at 750 ºC, 2 h to 33% at 900 ºC for 1 h. The high peaks at 60º correspond to the twin (or \( \Sigma 3 \)) GBs. It can be seen that at 600 ºC the fraction of \( \Sigma 3 \) GBs is below 30%, while between 750 ºC and 900 ºC it suddenly increases above 60%.

Figure 5. Cont.
4. Discussion

When one considers the conventional phase diagrams for three-dimensional phases, one does not expect any phase transformations in a single-phase region (like α solid solution). However, if one includes the grain boundaries (GBs) and their triple junctions (GB TJs), one can observe in GBs and GB TJs the thin layers (or thin channels) of the second phase. This is not stable in the bulk, and only one phase can exist (see, for example, References [8,45,68,114,148–150] and references therein). Most frequently the thin layers (or thin channels) of a metastable second phase are observed in GBs (or GB TJs) in the proximity of a solidus or solvus lines in a bulk phase diagram. Nevertheless, more recently one can find in the literature multiple indications of unusual phenomena which can take place removed from the solidus or solvus lines, in other words, “deep” in the solid-solution area. Thus, (i) an Al-based alloy can exhibit the anomalous superductility around 400 °C [165]; (ii) the additional
peaks appear in the DSC curves of the fine-grained Al-Mg alloys. These peaks were observed in a solid-solution region of the Al-Mg phase diagram [166,167]; and (iii) preliminary heat treatment in a single-phase region can modify thehardening and softening behavior of the UFG Al-Zn alloys during plastic deformation at room temperature [168]. This unusual behavior has been observed in UFG materials and, thus, can be due to the various GB phase transformations which could take place also in the one-phase area of a bulk phase diagram far away from solvus or solidus lines.

Let us start the discussion with Al-Zn alloys. In the Al-Zn system, numerous GB phase transitions have been observed. Historically, first, the transition between complete and incomplete GB wetting has been observed in the (Al) + L two-phase area of the Al-Zn phase diagram (where L means the liquid phase or melt and (Al) means the solid solution) [28]. In this two-phase area of the Al-Zn phase diagram, the (Al) solid solution exists in equilibrium with the melt. The completely wetted (Al)/(Al) GBs appear first above $T_{\text{wmin}} = 440^\circ$C. The respective tie-line at this temperature $T_{\text{wmin}}$ is shown in Figure 2a (thin solid line). The low index “min” means that this is the minimal temperature where the first (Al)/(Al) GBs completely wetted by the Al + Zn melt appear in the studied polycrystals. With increasing temperature, the amount of completely wetted (Al)/(Al) GBs continuously increases and reaches the maximum of 100% at $T_{\text{wmax}} = 565^\circ$C [28]. Above this temperature $T_{\text{wmax}}$, all GBs in the studied polycrystals are completely wetted by the liquid phase. Using DSC, one can observe the thermal effects of GB wetting phase transition [148]. It is important to note that this GB wetting phase transition in Al-Zn alloys is of the first order. Particularly, this fact reflects itself in temperature dependence of the GB contact angle $\theta$. The plot of the temperature dependence of contact angle between GB and melt is convex. Moreover, its first derivative has a discontinuity (or break) at $T_w$. In other words, the temperature dependence of $\theta$ has two different tangents at $T_w$. Nevertheless, the GB wetting phase transitions of the second order (or continuous transitions) also exist. They have been also observed in the Al-Zn system, but in the two-phase (Zn) + L area of the Al-Zn phase diagram [29]. In case of continuous GB wetting transition, the temperature dependence of $\theta$ is concave, and its first derivative has no discontinuity at $T_w$. In other words, the $\theta(T)$ dependence has only one tangent at $T_w$.

The GB wetting phase transitions are marked by the horizontal tie-lines in a two-phase “solid + liquid” area of a bulk phase diagram. These tie-lines should have a continuation in the one-phase solids-solution area of a phase diagram called GB solidus line. Such GB solidus also exists in the (Al) area of the Al-Zn phase diagram. It is shown in Figure 2a by the dotted line. It starts from the bulk solidus at $T_{\text{wmin}} = 440^\circ$C and finishes at the melting point of aluminum. Between bulk solidus and GB, the grain boundaries possess the thin layer of a liquid-like phase [68,88,169–175]. The GB solidus (also called the premelting transition line) has been observed experimentally by TEM [68]. In this work, the GB structure has been compared in polycrystalline samples quenched from the temperature above and below the bulk solidus. The GB premelting transition can also be observed using DSC. Thus, the GB melting in UFG or nanocrystalline samples starts earlier than in the bulk [68,114,176,177]. These data permit one to estimate the position of GB solidus in Al-Zn alloys.

The polycrystals containing a thin liquid-like layer in GBs become very ductile. For example, the unusual superplasticity with maximal elongation of up to 2500% has been observed in the Al-Zn-Mg UFG polycrystals [115–119]. This phenomenon appeared in the very narrow temperature interval below the solidus line. With increasing content of magnesium and zinc, the high superductility disappears [115–119]. Both superplasticity itself and the narrow (temperature and concentration) interval of its existence are unexplained. Only the observation of the GB wetting transition in Al-Zn, Al-Mg and in Al-Zn-Mg alloys gives the key to understanding this phenomenon [109–119,176,177]. Namely, when the GB wetting transitions take place in the two-phase (Al,Mg,Zn) + L area of the phase diagram, then the respective wetting tie-lines continue the into the (Al,Mg,Zn) solid–solution area with GB solidus lines. Therefore, between bulk and GB solidus, the grain boundaries contain thin liquid-like layers which ensure the high plasticity [109–114,178].
Another class of GB phase transitions is also observed in the Al-Zn alloys. In this case a GB can be completely (or partially) wetted by a second solid phase instead of a liquid one. The thermodynamics of this wetting transition are similar to that for the liquid phase. In other words, a continuous layer of a second phase (no matter, liquid or solid) substitutes a grain boundary if the GB energy $\sigma_{GB}$ is higher than the energy $2\sigma_{SL}$ of two solid/liquid boundaries or the energy $2\sigma_{\alpha\beta}$ of two boundaries between two solid phases $\alpha$ and $\beta$. The GB wetting transition by a second solid phase has been observed for the first time in the Al-Zn system [179,180]. It takes place in the (Al) + (Zn) two-phase area of the Al-Zn binary phase diagram where the Al-rich (Zn) solid solution is in equilibrium with the Zn-rich (Al) solid solution. It is just below the (Al) + (Zn) eutectic point at $T_e = 381^\circ$C (see Figure 2a). The important difference between GB wetting by a liquid phase and by a second solid phase is that the kinetics of a latter are very slow in comparison with the first. For example, one needs to wait hundreds of hours to observe how the continuous (Al) GB layer between two (Zn) grains transforms into a chain of lenticular particles instead of the few minutes needed to equilibrate liquid phase layers between solid grains. The percentage of boundaries between (Zn) grains completely wetted by the (Al) face-centered cubic (fcc) solid phase increases from zero at $T_{wminZn} = 283 ^\circ$C up to about 30% at eutectic point $T_e = 381 ^\circ$C.

A similar transformation in (Al)/(Al) GBs proceeds in the (Al') + (Zn) area of the Al-Zn binary phase diagram below the temperature of monotectoid decomposition $T_{mon} = 277 ^\circ$C (see Figure 2a) [150,179,181]. In this case, the (Al')/(Al') GBs between grains of (Al') solid solution with low zinc content could be completely wetted by a layer of the second solid phase (Zn). Differently from the previous case, the percentage of such GBs increases with decreasing temperature. The first (Al')/(Al') GBs completely wetted by a continuous layer of the second solid phase (Zn) appear below $T_{wmin} = 205 ^\circ$C. The portion of (Al)/(Al) GBs completely wetted by the second solid phase (Zn) increases with decreasing temperature. It becomes 100% at nearly $T_{wmax} = 125 ^\circ$C. Unfortunately, it is not easy to measure $T_{wmax}$ because the equilibration of a microstructure runs extremely slowly below 200 $^\circ$C. However, it is possible to apply HPT to accelerate the mass-transfer in the (Al') + (Zn) alloys [182,183]. This is possible due to the high amount of non-equilibrium vacancies appearing during HPT. After HPT of Al-Zn alloys, it was possible to observe the so-called pseudoincomplete (or pseudopartial) GB wetting [113,184,185]. In this case, the few nanometer thin Zn-rich layer exists in the (Al)/(Al) GBs parallel with a non-zero contact angle between (Zn) particles and GBs. After HPT treatment of Al-Zn alloys, about one-third of (Al)/(Al) GBs contain such thin Zn-rich layer [113]. We discussed above that the thin liquid-like GB layers close to the bulk solidus lines are the reason for the superplasticity of UFG Al-based alloys [109–119]. Similarly, the thin solid Zn-rich layers in the solid Al-Zn UFG polycrystals appeared after HPT leads to their high ductility [109–114,186]. The pseudoincomplete GB wetting in the Al-Zn polycrystals is not surprising. For example, the GB wetting in the Al-Zn alloys can be continuous or discontinuous [29]. The theory predicts that the line of the discontinuous GB transitions from complete to incomplete wetting can divide into two lines; one of first and one of second order transformation. In this case, the area of pseudoincomplete GB wetting should arise between these two lines [187–199].

It is clear that the tie-lines of GB phase transformations existing in the (Al) + (Zn) two-phase area of the Al-Zn binary phase diagram could have a continuation into one-phase solid–solution (Al) area. By crossing of these lines, the peaks in DSC curves would appear. The small endothermic peaks observed in DSC curves (Figure 3a) could indicate the presence of such GB lines. We have drawn them between the bulk solvus line and two lower experimental DSC as crosses in Figure 3a. By drawing such GB lines in the bulk phase diagram, we have to remember some important circumstances. When the first evidence of the presence of GB phase transformations appeared, these experimental findings are discussed in terms of transitions “pure (or dry) GB → few nanometer thin GBF → thick GB layer of bulk phase”. Later, it has been shown that the variety of possible GB states (called complexions) is much richer (see for example References [17,19–27,87,169,170]). One can imagine the transitions from clean GB to single-layer GBF, then to bi-layer, tri-layer etc. [17,19–27] Some of these transition lines can
We can suppose that the formation of other (less thick) GB complexions can take place below \( T \). Above \( (\text{dotted}) \) tie-line is drawn in Figure 2a in the \((\text{Al}') + (\text{Al}'')\) area. According to Cahn’s prediction, \( (\text{Figure 3b}) \) is below the \( (\text{Figure 2b}) \) [148,176,201,202]. The respective tie-lines at \( T_{\text{min}} = 540 \) °C and \( T_{\text{max}} = 610 \) °C are shown in the Al-Mg phase diagram (Figure 2b). The upper endothermic peak observed in the DSC curves (Figure 3b) is below the \( T_{\text{min}} = 540 \) °C. However, the tie-line at \( T_{\text{min}} = 540 \) °C marks the formation of macroscopically thick continuous GB layer of a liquid phase. Following the discussion in Reference [20], we can suppose that the formation of other (less thick) GB complexions can take place below \( T_{\text{min}} \). Therefore, we may draw another tie-line in the \((\text{Al}) + L\) area below \( T_{\text{min}} = 540 \) °C and construct its continuation into the one-phase area of solid solution \((\text{Al})\) (see Figure 3b). Such continuation can also finish at the bulk solidus (as in References [37,39,42,173–175] or with a critical point as in References [17,19,20,23,27,87]).

The conventional wetting of \((\text{Al})/(\text{Al})\) GBs by the melt was also observed in the Al-Mg alloys (Figure 2b) [148,176,201,202]. The respective tie-lines at \( T_{\text{min}} = 540 \) °C and \( T_{\text{max}} = 610 \) °C are shown in the Al-Mg phase diagram (Figure 2b). The upper endothermic peak observed in the DSC curves (Figure 3b) is below the \( T_{\text{min}} = 540 \) °C. However, the tie-line at \( T_{\text{min}} = 540 \) °C marks the formation of macroscopically thick continuous GB layer of a liquid phase. Following the discussion in Reference [20], we can suppose that the formation of other (less thick) GB complexions can take place below \( T_{\text{min}} \). Therefore, we may draw another tie-line in the \((\text{Al}) + L\) area below \( T_{\text{min}} = 540 \) °C and construct its continuation into the one-phase area of solid solution \((\text{Al})\) (see Figure 3b). Such continuation can also finish at the bulk solidus (as in References [37,39,42,173–175] or with a critical point as in References [17,19,20,23,27,87]).

In the Al-Zn binary phase diagram, a very interesting two-phase area exists between \( T_{\text{mon}} = 277 \) °C and \( T_{c} = 351.5 \) °C (see Figure 2a). Below the critical point at \( T_{c} = 351.5 \) °C, the single fcc \((\text{Al})\) solid solution decomposes into two fcc solid solutions, namely a Zn-poor (\( \text{Al}' \)) and a Zn-rich (\( \text{Al}'' \)) one. The seminal work of Cahn [1] discusses such phase diagrams with critical points. Cahn noticed that in the critical point the difference between the two phases, as (\( \text{Al}' \)) and (\( \text{Al}'' \)) in our case, disappears. Therefore, the energy of interphase boundary \( 2\gamma_{\text{Al},\text{Al}''} \) tends to zero by approaching to \( T_{c} \) from below. It means that the phase (\( \text{Al}'' \)) should wet the \((\text{Al}')/(\text{Al}')\) GB close enough to \( T_{c} \). Such a hypothetical (dotted) tie-line is drawn in Figure 2a in the \((\text{Al}') + (\text{Al}'')\) area. According to Cahn’s prediction, the tie-line in the \((\text{Al}') + (\text{Al}'')\) area should continue in the \((\text{Al}')\) one-phase area and finish in a critical point [1]. We have drawn such a (dotted) line through the upper cross which marks the upper small endothermic peak observed in DSC curves (Figure 3a). However, it is not the only unique possible construction to explain the upper point. The Al-Zn system opens many possibilities to construct the lines for GB complexions transitions and may also finish in a critical point [17,19–27,87,169,170].

Moreover, it is known [203] that the bulk phase diagram of Al-Zn can be more complex than the diagram presented in the handbook [161]. Thus, Krishna Rao et al. observed a metastable phase in the Al-50 wt.% Zn alloy in the single-phase (Al) region at 400 °C [204]. This phase had a trigonal structure with space group \( R3m \) and unit cell period \( a = 0.2852 \text{ nm} \) and \( c = 0.685 \text{ nm} \). In addition, A. Sandoval-Jimenez et al. observed the phase with a triclinic structure in the Al-78.2 wt.% Zn alloy [205]. This means that in a single-phase region (Al), the metastable phases can exist. In the Al-Mg system there were also observed several metastable phases [167,203,206,207].

The Cu-Ag phase diagram (Figure 2e) is a eutectic one similar to the Al-Zn phase diagram. The \((\text{Cu})/(\text{Cu})\) GBs can be either completely or incompletely wetted by the Ag-rich liquid phase. The first high-angle \((\text{Cu})/(\text{Cu})\) GBs completely wetted by the Ag-rich melt appear above \( T_{\text{wmin}} = 790 \) °C [208]. The \( T_{\text{wmin}} = 790 \) °C is very close to the Cu-Ag eutectic temperature \( T_{e} = 779 \) °C [161]. Above \( T_{\text{wmax}} = 1020 \) °C, all high-angle \((\text{Cu})/(\text{Cu})\) GBs are completely wetted by the Ag-rich melt [208]. Later, it has been found that above \( T_{\text{wmax}} = 1020 \) °C, the low-angle GBs (with the energy below that of the high-angle GBs) start to be completely wetted by the liquid phase [208]. The tie-lines of GB wetting
phase transitions (thin solid lines) in the (Cu) + L two-phase area of the Cu-Ag phase diagram are shown in Figure 2e.

Similar to the Al-Zn and Al-Mg phase diagram, these tie-lines of GB wetting phase transitions could continue in the (Ag) one-phase area as the lines of GB premelting. Indeed, the DSC curves for UFG Cu-2.8 wt.% Ag, Cu-4 wt.% Ag, Cu-5 wt.% Ag, and Cu-6.5 wt.% Ag alloys contain additional peaks below the bulk solidus, similar to the Al-Zn and Al-Zn-Mg alloys [68,109–114,176–178]. Moreover, we can determine in each DSC curve for UFG samples one strong (open circles) and possibly one weak peak (open squares) below the bulk solidus (full down-triangles). This could be an indication of different kinds of GB complexions below the bulk solidus line. The DSC curves for the second heating, i.e., for coarse-grained samples, are more simple in comparison with DSC curves for the UFG samples. We can see that in case of the second heating, only the DSC curves for Cu-5 wt.% Ag and Cu-6.5 wt.% Ag alloys have a complicated melting peak and the curves for the other two alloys (with less silver) do not. The tangents in the main melting peak correspond to the liquidus and solidus temperature and coincide well with each other for the first and second heating with only a few degrees difference. The respective points are shown in the Cu-Ag phase diagram (Figure 2e) with up- and down- full triangles. The additional melting peaks are more pronounced in the UFG samples and lie at about 10–30 °C below the bulk solidus line and are shown in the Cu-Ag phase diagram (Figure 2e) by open circles and squares. The difference between bulk and GB solidus increases with increasing silver content from ~10 °C to ~30 °C (Figures 3d and 2e). The peaks observed in the DSC curves (Figure 3d,e) close to 778 °C are shown by crosses. Open diamond shows that such a peak is not present at 2.8 wt.% Ag. Seven hundred and seventy-eight degrees Celsius is very close to $T_e = 779$ °C, but, nevertheless, is slightly below it. What is the possible reason for these peaks?

The GB segregation in Cu-Ag alloys has been analyzed in several papers on numerical computation. In particular, it has been predicted the transition from monolayer to multilayer, silver segregation can take place in (Cu)/(Cu) GBs, or even to the thick premelting Ag layer [209–212]. In Reference [72], a series of discontinuous GB layering transitions has been numerically investigated for the Cu-Ag alloys. They are associated with GB segregation in a lattice–gas model of a binary Cu-Ag alloy. The Monte-Carlo simulations of the [001]-axis symmetric tilt $\Sigma5$ (310) GB in Cu-Ag alloys have been performed in References [210–212]. In References [213,214], the effective Ising model (mean-field theory) has been applied to the Cu-Ag system. These works predicted the first order monolayer-to-multilayer phase transition in the first layers near the GB plane. Such GB phase transition can be considered as a generalized Fowler-Guggenheim one. It can flow by the formation of a thick prewetting GB layer [215]. The role of Bi, Na, and Ag as substitutional impurities in the $\Sigma5$ (310) [001] symmetric tilt (Cu)/(Cu) GB has been analysed [216]. These first-principles calculations explained why bismuth can cause the embrittlement of (Cu)/(Cu) GBs and silver cannot.

Thus, we may suppose that similar to the Al-Zn and Al-Mg systems, a kind of GB complexion transition takes place at 778 °C in the one phase (Cu) area of the Cu-Ag phase diagram (Figure 2e). This may be the transition from one-layer to the multilayer GB segregation. We mark this transition in the Cu-Ag phase diagram (Figure 2e) by the thin dotted line. However, the observed thermal effect of this GB transition decreases with decreasing Ag content and becomes completely invisible at 2.8 wt.% Ag (Figure 3d). Therefore, the dotted line in the (Cu) area finishes with a critical point between 2 and 4 wt.% Ag. In the two-phase (Cu) + (Ag) area of the Cu-Ag phase diagram, the composition of (Cu) and (Ag) solid solution is constant at constant temperature. Therefore, the dotted line of GB complexion transition continues in the two-phase (Cu) + (Ag) area as a horizontal tie-line at 778 °C.

The Cu-Co phase diagram is a peritectic transition in the Cu-rich side. In this case, Co content in the solid solution is higher than that in the liquid phase as long as we consider the two-phase solid + liquid area of the Cu-Co phase diagram. However, it has been observed that this cobalt-poor melt can completely wet the (Cu)/(Cu) GBs [32]. We constructed the GB solidus line (or GB premelting line) just below the bulk solidus line (dotted line in Figure 2) using the DSC data [177]. The new data on GB DSC peaks are marked by two crosses. They lie well below the GB solidus line. We would
explain them by the (not observed before) GB complexions transition. The respective hypothetical line of such transition is also drawn in Figure 2d. Zhevnenko et al. studied the creep of the Cu-Co solid solutions close to the bulk solidus [164,217–219]. Zhevnenko et al. found a sudden break in the temperature dependences of creep rate even lower than the DSC points shown by crosses in Figure 2d. These breaks can also be due to the GB complexions transition.

In Figure 4c, dislocations are visible in the TEM. Actually, the dislocations can also enhance the area for precipitation and provide alterations to the wetting conditions. They can, themselves, be incompletely or completely wetted by the second phase. In the first case, a chain of droplets (or particles) form along the dislocation. In the latter case, the continuous “tube” of a melt or a second solid phase can form along the dislocation. If the dislocations can build low-angle GBs, the temperature of wetting transition for them differs from that for the high-angle GBs [208]. From the wetting point of view, there is some similarity between dislocations and another line defect—namely GB triple junctions [30].

The presence of GBFs (or more general—GB complexions) strongly influences the grain growth as well as GB character distribution in oxides and metallic alloys [54,56–59,63,75,76,79,80,83,106,123–130,151–160]. It has been known for a long time that impurities can shift the borders between GBs with special and general structure and properties [142,143,220,221]. The GB complexions enriched by the impurities form, most likely, in the GBs with high energy (also called general GBs) [79] and, less likely, in coincidence GBs with dense and nearly perfect structure [123]. During grain growth, the competition between various GBs with different mobilities takes place [23,151–160]. Moreover, the micrographs in Figure 5a,c permit to observe that the distribution of small and large grains is not uniform. It indicates that the grain growth in the sample is abnormal, at least, in some locations. Therefore, if we have the lines of GB complexion transition in a bulk phase diagram, the portion of most perfect GBs Σ3 (or twin GBs) can serve as an indirect indication for such GB phase transformations. The Cu-Ni system is quite promising for the observation of these GB complexion transitions. The Cu-Ni phase diagram (Figure 2d) contains at high temperatures, the continuous raw of fcc solid solution. However, below the critical point, \( T_c = 354.5 \, ^\circ C \), these solid solutions decompose into Cu-rich and Ni-rich ones. Similar to the Al-Zn system, following Cahn’s idea [1], one can expect that under the decomposition dome, close enough to the critical point \( T_c \) the phase (Cu) should wet the (Ni)/(Ni) GBs. Such a hypothetical (dotted) GB wetting tie-line is drawn in Figure 2d in the (Cu) + (Ni) area. Again, following the Cahn’s prediction, the GB wetting tie-line in the (Cu) + (Ni) area should continue into the (Ni,Cu) one-phase area and finish in a critical point [1]. We have drawn such a (dotted) line through the middle experimental point at 750 \( ^\circ C \), since above this temperature the portion of Σ3 GBs in a Cu-77 wt. % Ni polycrystal suddenly increases from 0.3–0.37 at 600–750 \( ^\circ C \) to 0.6 at 900 \( ^\circ C \). Moreover, some microtwins are visible even in UFG samples after HPT (shown by arrows in Figure 1c). Thus, the changes in GB character distribution during grain growth, as well as a certain input of abnormal grain growth, can serve as an indication of GB complexions transitions [17,19–27,87,169,170]. Of course, this phenomenon is only an indirect indication of possible GB complexions transition because the increase of the portion of less mobile low-Σ GBs during grain growth cannot be called unexpected by temperature increase above 0.5 \( T_m \). However, the observed effect is very promising for a new understanding of GB complexions and permits to plan the more detailed studies of GB structure and properties in this system.

It is clear that the new GB lines in the Al-Zn, Al-Mg, Cu-Co, and Cu-Ni phase diagrams proposed in this work are still very hypothetic, and the physical reasons for the small endothermic peaks observed in DSC curves (Figure 3) should be studied in more details.

5. Summary and Conclusions

Small endothermic peaks were observed in ultra-fine grained (UFG) Al-Zn, Al-Mg, Cu-Ag, and Cu-Co solid solutions during their heating in the differential scanning calorimeter (DSC). These peaks are far away from the solidus and solvus lines of respective bulk phase diagrams.
The endothermic effect is more pronounced in UFG samples with a high specific area of grain boundaries (GBs) than in their coarse-grained counterparts. In Cu-Ni UFG solid solution the amount of twin GBs drastically increases with increasing temperature. We explain these observed phenomena by the phase transformations between GB complexions. The new lines of these GB phase transformations have been constructed in the Al-Zn, Al-Mg, Cu-Ag, Cu-Co, and Cu-Ni bulk phase diagrams.

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